Interpreting Xor Intuitionistic Fuzzy Connectives from **Quantum Fuzzy Computing**

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Computer systems based on intuitionistic fuzzy logic are capable of generating a reliable output even when Abstract: handling inaccurate input data by applying a rule based system, even with rules that are generated with imprecision. The main contribution of this paper is to show that quantum computing can be used to extend the class of intuitionistic fuzzy sets with respect to representing intuitionistic fuzzy Xor operators. This paper describes a multi-dimensional quantum register using aggregations operators such as t-(co)norms based on quantum gates allowing the modeling and interpretation of intuitionistic fuzzy Xor operations.

INTRODUCTION 1

Fuzzy logic (FL) and its extensions as the Atanassov's fuzzy intuitionistic logic (A-IFL)(Atanassov and Gargov, 1998; Atanassov, 2017) together with quantum computing (QC) are relevant research areas consolidating the analysis and the search for new solutions for difficult problems faster than the classical logical approach by extending results from conventional computing.

Similarities between these areas in the representation and modelling of uncertainty have been explored (Pykacz, 1993; Kreinovich et al., 2009; Mannucci, 2006). The uncertainty of human being's reasoning can be modelled in A-IFL as a mathematical model inheriting the indeterminacy from membership and non-membership degrees which are not necessarily complementary. It provides techniques helping physicists and mathematicians to transform their uncertainty ideas into new computational programs (Kosheleva et al., 2015).

The uncertainty of the real world is concerned with fundamental concepts of QC by making use of properties of quantum mechanics as the superposition which suggest an improvement in the efficiency regarding complex tasks. In addition, simulations using classical computers improve the development and validation of basic quantum algorithms, anticipating the knowledge related to their behaviours when executed in a quantum computer.

In spite of quantum computers being restricted to a few research centers and laboratories, the studies from quantum information and quantum computation are a reality nowadays.

In this context, new methods dealing with fuzzy approaches to quantum computational logics have been proposed (Chiara et al., 2018). Moreover, it contributes to increase the interest in quantum algorithm applications representing fuzzy and intuitionistic fuzzy systems.

1.1 Synergy between FL and QC

Potentialities from quantum parallelism and superposition of quantum states are explored by new methodologies. See, e.g. providing a model for humanoid behaviours based on FL (Raghuvanshi and Perkowski, 2010) and using a density operator and logical connectives defined by quantum gates (Bertini and Leporini, 2017). In other areas, QC is used to improve computations as Econometric Modeling (Sriboonchitta et al., 2019) and to integrate concepts from quantum genetic algorithm and fuzzy neural network (Li et al., 2010) as a model employed to control an inverted pendulum system.

From a logical perspective, many-valued fuzzy approach is a consolidated research area (Hájek, 1998) and recent efforts integrating the quantum computational logic are considered (Chiara et al., 2018). See, e.g. (Freytes et al., 2010) making use of quantum states to represent information processing in QC.

In (Rigatos and Tzafestas, 2002), a parallel control fuzzy algorithm is proposed showing that both QC and fuzzy learning algorithms are relevant research

288

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areas which can profit from each other.

In this work, the modelling and interpretation of A-IFL via QC provides the description of intuitionistic fuzzy connectives by using quantum registers and quantum transformations from the traditional model of quantum circuits.

The information regarding each intuitionistic fuzzy set is represented by pairs of quantum registers, guaranteeing the inherent unitarity of quantum states and quantum transformations well as the flexibility of the complementarity relation of membership and non-membership functions characterizing intuitionistic fuzzy sets.

Some results of such research area are mainly related to interpretation of fuzzy connectives via QC, as negation, conjunction, disjunction, implications and xor operators (Avila et al., 2015) and to intuitionistic fuzzy connectives (Reiser et al., 2016).

This work extends that approach by the interpretation of two classes of intuitionistic fuzzy Xor connectives via QC, the \otimes_I and \oplus_I intuitionistic fuzzy Xor operators, which can be expressed as composition of conjunctions (S_I), disjunctions (T_I) and the complementary operators (N_I). But different from other representable intuitionistic fuzzy connectives, an intuitionistic fuzzy Xor connective cannot be expressed as an *N*-dual pair of Xor and EXor fuzzy connectives.

So, quantum states and quantum operators provide interpretation for intuitinistic fuzzy values and connectives, respectively. In particular, Xor and EXor can be applied to explore regular applications based on symmetric functions by making use of reversible logic (He et al., 2017) and can be also extended to QC (Perkowski et al., 2001).

1.2 Paper Outline

The remainder of this paper is organized as follows. Section 2 presents the foundations on A-IFL. Section 3 brings the main concepts of QC. In Section 4, the study and modeling of intuitionistic fuzzy Xors using QC is described. An interpretation of classical A-IFS (Atanassov, 2016; Mannucci, 2006) from quantum states is also considered, presenting the operations on A-IFS modelled from quantum transformations, relating the Atanassov's intuitionistic fuzzy approach(Atanassov, 1986) to two classes of Xor operators and also considering representable Xor connectives obtained by composition of standard negation together with the Product and Algebraic Sum. Finally, conclusions and further work are discussed in Section 5.

2 FUZZY LOGIC APPROACH

The A-IFL is a type-2 fuzzy logic conceived as a generalization of FL overcoming the limitations related to fuzzy sets for dealing with problems where the rules applied to the system could not be defined with precision, mainly related to non-membership degree which cannot be defined as a complement of its membership degree.

An element $x \in X$ belongs to the subset *A* such that $0 \le \mu_A(x) + \nu_A \le 1$, meaning that a non-membership degree $\nu_A(x)$ is not necessary the complement of its membership degree $\mu_A(x)$. Thus, *A* is given as:

$$A = \{(x, \mu_A(x), \nu_A(x)) : x \in \mathcal{X}, 0 \le \mu_A(x) + \nu_A \le 1\}$$
(1)

Taking $\mu_A(x) = x_1, \nu_A(x) = x_2 \in [0, 1]$ for an element $x \in \chi$, the set of all Atanassov's intuitionistic fuzzy values is given as $\tilde{U} = \{(x_1, x_2) : x_1 + x_2 \le 1\}$. The least and greatest element on \tilde{U} are given as $\tilde{0} = (0, 1)$ and $\tilde{1} = (1, 0)$, respetively.

Considering the research of A-IFL, it makes possible to extend the usual logic connectives, as follows:

- 1. *Conjunction*, usually modelled by a *triangular norm* (t-norm) operator (Klement et al., 2013), which is an additive aggregation used in the framework of intersection in A-IFL;
- 2. *Disjunction* which is frequently modelled by a *triangular conorm* (Klement et al., 2013) (t-conorm) representing the t-norm dual operation;
- 3. *Complement* as a negation, a non-increasing function, defined in the Atanassov's seminal work (Atanassov, 2006), reverting the extremes of the unit interval [0, 1].
- 4. *Xor operators*, considering in this work two representable classes \oplus_{I_P} and \otimes_{I_P} (Bedregal et al., 2013), both defined by composition of the above connectives verifying the symmetry, associativity, neutral element and boundary conditions.

Now, the connectives used to make the correlation between QC and A-IFL will be described in terms of representability based on fuzzy negations and aggregations. The standard intuitionistic fuzzy negation (Atanassov, 2006) is expressed as follows:

$$N_{I_{S}}(\tilde{x}) = (x_{2}, x_{1}), \forall \tilde{x} = (x_{1}, x_{2}) \in \tilde{U}.$$
 (2)

Meanwhile, for all $\tilde{x} = (x_1, x_2), \tilde{y} = (y_1, y_2) \in \tilde{U}$, the intersection and union can be defined, respectively, in terms of a t-norm *T* and a t-conorm *S*, given as:

 $T_{I}(\tilde{x}, \tilde{y}) = T_{I}((x_{1}, x_{2}), (y_{1}, y_{2})) = (T(x_{1}, y_{1}), S(x_{2}, y_{2}));$ $S_{I}(\tilde{x}, \tilde{y}) = S_{I}((x_{1}, x_{2}), (y_{1}, y_{2})) = (S(x_{1}, y_{1}), T(x_{2}, y_{2})).$ We consider the Product t-norm $T_{I_{P}}$ and Algebraic Sum $S_{I_{P}}$, respectively described as follows:

 $T_{I_P}(\tilde{x}, \tilde{y}) = (T_P(x_1, y_1), S_P(x_2, y_2)) = (x_1y_1, x_2 + y_2 + x_2y_2);$ $S_{I_P}(\tilde{x}, \tilde{y}) = (S_P(x_1, y_1), T_P(x_2, y_2)) = (x_1 + y_1 + x_1y_1, x_2y_2)).$ The classical xor expressions considered for this work are described as follows:

$$\begin{array}{lll} A \oplus B & \equiv & (\neg A \wedge B) \lor (A \wedge \neg B); \\ A \otimes B & \equiv & (A \lor B) \land (\neg A \lor \neg B). \end{array}$$

Definition 2.1. The Atanassov's intuitionistic fuzzy Xor operator is a function $E_I : \tilde{U}^2 \to \tilde{U}$ verifying:

 $\begin{aligned} EI: \ E_{I}(\tilde{0},\tilde{0}) &= E_{I}(\tilde{1},\tilde{1}) = \tilde{0} \text{ and } E_{I}(\tilde{0},\tilde{1}) = E_{I}(\tilde{1},\tilde{0}) = \tilde{1}; \\ E2: \ E_{I}(\tilde{x},\tilde{y}) &= E_{I}(\tilde{y},\tilde{x}); \\ E3: \ \tilde{y}_{1} \leq \tilde{y}_{2} \Rightarrow E_{I}(\tilde{0},\tilde{y}_{1}) \leq E_{I}(\tilde{0},\tilde{y}_{2}); \\ E4: \ \tilde{y}_{1} \leq \tilde{y}_{2} \Rightarrow E_{I}(\tilde{1},\tilde{y}_{1}) \geq E_{I}(\tilde{1},\tilde{y}_{2}). \end{aligned}$

Proposition 2.1. The functions $\oplus_I, \otimes_I : \tilde{U}^2 \to \tilde{U}$ given as follows:

are intuitionistic fuzzy Xor operators.

Proof. Consider \oplus_I operator. It holds that:

 E_I 1: The following boundary conditions are verified in the endpoints of unit interval:

 $\begin{array}{l} \oplus_{I}(\tilde{0},\tilde{0}) = S_{I}(T_{I}(N_{SI}(\tilde{0}),\tilde{0}),T_{I}(\tilde{0},N_{SI}(\tilde{0}))) = S_{I}(\tilde{0},\tilde{0}) = \tilde{0}; \\ \oplus_{I}(\tilde{1},\tilde{1}) = S_{I}(T_{I}(N_{SI}(\tilde{1}),\tilde{1}),T_{I}(\tilde{1},N_{SI}(\tilde{1}))) = S_{I}(\tilde{0},\tilde{0}) = \tilde{0}; \\ \oplus_{I}(\tilde{0},\tilde{1}) = S_{I}(T_{I}(N_{SI}(\tilde{0}),\tilde{1}),T_{I}(\tilde{0},N_{SI}(\tilde{1}))) = S_{I}(\tilde{1},\tilde{0}) = \tilde{1}; \\ \oplus_{I}(\tilde{1},\tilde{0}) = S_{I}(T_{I}(N_{SI}(\tilde{1}),\tilde{0}),T_{I}(\tilde{1},N_{SI}(\tilde{0}))) = S_{I}(\tilde{0},\tilde{1}) = \tilde{1}; \end{array}$

 $E_I 2: \forall \tilde{x}, \tilde{y} \in \tilde{U}$ the following is verified

$$E_I(\tilde{x}, \tilde{y}) = S_I(T_I(N_{SI}(\tilde{x}), \tilde{y}), T_I(\tilde{x}, N_{SI}(\tilde{y})))$$

= $S_I(T_I(\tilde{y}, N_{SI}(\tilde{x})), T_I(N_{SI}(\tilde{y}), \tilde{x})) = E_I(\tilde{y}, \tilde{x});$

 E_I 3 : If $\tilde{y}_1 \leq \tilde{y}_2$ then the following is verified

$$\begin{split} \oplus_{I}(\tilde{0},\tilde{y}_{1}) &= S_{I}(T_{I}(N_{SI}(\tilde{0}),\tilde{y}_{1}),T_{I}(\tilde{0},N_{SI}(\tilde{y}_{1}))) \\ &\leq S_{I}(T_{I}(N_{SI}(\tilde{0}),\tilde{y}_{2}),T_{I}(\tilde{0},N_{SI}(\tilde{y}_{2}))) = E_{I}(\tilde{0},\tilde{y}_{2}); \end{split}$$

$$\begin{split} E_{I}4: & \text{If } \tilde{y}_{1} \leq \tilde{y}_{2} \text{ then the following is verified} \\ \oplus_{I}(\tilde{1}, \tilde{y}_{1}) = S_{I}(T_{I}(N_{SI}(\tilde{1}), \tilde{y}_{1}), T_{I}(\tilde{1}, N_{SI}(\tilde{y}_{1}))) \\ \geq S_{I}(T_{I}(N_{SI}(\tilde{1}), \tilde{y}_{2}), T_{I}(\tilde{1}, N_{SI}(\tilde{y}_{2}))) = E_{I}(\tilde{1}, \tilde{y}_{2}); \end{split}$$

Therefore, Proposition 2.1 is verified.

Based on those expressions, one can use the operators T_I , S_I , N_{I_S} to construct the Atanassov's intuitionistic fuzzy xor (\oplus_I, \otimes_I) respectively given as follows:

$$\begin{split} &\oplus_{I}(\tilde{x},\tilde{y}) = \\ &= S_{I}(T_{I}(N_{SI}(\tilde{x}),\tilde{y}),T_{I}(\tilde{x},N_{SI}(\tilde{y}))) \\ &= S_{I}(T_{I}((x_{2},x_{1}),(y_{1},y_{2})),T_{I}((x_{1},x_{2}),(y_{2},y_{1})))) \\ &= S_{I}((T(x_{2},y_{1}),S(x_{1},y_{2})),(T(x_{1},y_{2}),S(x_{2},y_{1})))) \\ &= (S(T(x_{2},y_{1}),T(x_{1},y_{2})),T(S(x_{1},y_{2}),S(x_{2},y_{1}))) (5) \\ &\otimes_{I}(\tilde{x},\tilde{y}) = \\ &= T_{I}(S_{I}(\tilde{x},\tilde{y}),S_{I}(N_{SI}(\tilde{x}),N_{SI}(\tilde{y}))) \\ &= T_{I}(S_{I}((x_{1},x_{2}),(y_{1},y_{2})),S_{I}((x_{2},x_{1}),(y_{2},y_{1}))) \\ &= (T(S(x_{1},y_{1}),T(x_{2},y_{2})),(S(x_{2},y_{2}),T(x_{1},y_{1}))) \\ &= (T(S(x_{1},y_{1}),S(x_{2},y_{2})),S(T(x_{2},y_{2}),T(x_{1},y_{1}))) (6) \end{split}$$

By considering T_{I_P} and S_{I_P} in Eqs. (5) and (6), the xor $(\bigoplus_I, \bigotimes_I)$ operators can be respectively expressed as follows:

$$\begin{split} \oplus_{I}(\tilde{x}, \tilde{y}) &= ((x_{1}y_{2} + x_{2}y_{1} - x_{1}x_{2}y_{1}y_{2}), \\ &= (x_{1}x_{2} + x_{1}y_{1} + x_{2}y_{2} + y_{1}y_{2} - x_{1}x_{2}y_{1} - x_{1}x_{2}y_{2} - x_{1}y_{1}y_{2} - x_{2}y_{1}y_{2} + x_{1}x_{2}y_{1}y_{2})). \end{split}$$
(7)
$$\otimes_{I}(\tilde{x}, \tilde{y}) &= ((x_{1}x_{2} + x_{1}y_{2} + x_{2}y_{1} + y_{1}y_{2} - x_{1}x_{2}y_{1} + x_{1}x_{2}y_{2}y_{1} + x_{1}x_{2}y_{1}y_{2}), \\ &= (x_{1}y_{1} + x_{2}y_{2} - x_{1}y_{1}y_{2} - x_{2}y_{1}y_{2} + x_{1}x_{2}y_{1}y_{2}), \\ &= (x_{1}y_{1} + x_{2}y_{2} - x_{1}x_{2}y_{1}y_{2})). \end{aligned}$$
(8)

According to (Mannucci, 2006), fuzzy sets can be obtained by quantum superposition of classical fuzzy states associated with a quantum register.

3 QUANTUM COMPUTING

In *QC*, the *qubit* is the basic information unit, being the simplest quantum system, defined by a unitary and bi-dimensional state vector.

Qubits are generally described, in Dirac's notation (Nielsen and Chuang, 2003), by the following expression

$$|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$$

when the coefficients α and β are complex numbers for the amplitudes of the corresponding states in the computational basis (state space), respecting the condition $|\alpha|^2 + |\beta|^2 = 1$, which guarantees the unitary of the state vectors of the quantum system, represented by $(\alpha, \beta)^t$ (Kave et al., 2007).

by $(\alpha, \beta)^t$ (Kaye et al., 2007). The state space of a quantum system with multiple *qubits* is obtained by the tensor product of the space states of its subsystems. Considering a quantum system with two *qubits*, $|\Psi\rangle = \alpha|0\rangle + \beta|1\rangle$ and $|\varphi\rangle = \gamma|0\rangle + \delta|1\rangle$, the state space comprehends the tensor product given by

$$|\Pi\rangle = |\psi\rangle \otimes |\phi\rangle = \alpha \cdot \gamma |00\rangle + \alpha \cdot \delta |01\rangle + \beta \cdot \gamma |10\rangle + \beta \cdot \delta |11\rangle.$$

The state transition of a quantum systems is performed by controlled and unitary transformations associated with orthogonal matrices of order 2^N , with *N* being the number of *qubits* within the system, preserving norms, and thus, probability amplitudes (Imre and Balázs, 2005). For instance, the *NOT* operator (*Pauli-X* transformation) and its application over 1-dimensional and 2-dimensional quantum systems are presented in the following.

$$X|\psi\rangle = \begin{pmatrix} 0 & 1\\ 1 & 0 \end{pmatrix} \begin{pmatrix} \alpha\\ \beta \end{pmatrix} = \begin{pmatrix} \beta\\ \alpha \end{pmatrix}; \quad (9)$$

$$X^{\otimes 2}|\Pi\rangle = \begin{pmatrix} 0 & 0 & 0 & 1\\ 0 & 0 & 1 & 0\\ 0 & 1 & 0 & 0\\ 1 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \alpha \cdot \gamma \\ \alpha \cdot \delta \\ \beta \cdot \gamma \\ \beta \cdot \delta \end{pmatrix} = \begin{pmatrix} \alpha \cdot \gamma \\ \alpha \cdot \delta \\ \beta \cdot \delta \\ \beta \cdot \gamma \end{pmatrix}$$
(10)

Furthermore, the action of a Toffoli QT is also shown in next Eq. (11), describing a controlled operation for a 3-dimensional quantum system

$$T|\chi\rangle = T(\psi \otimes \varphi \otimes \sigma).$$

In this case, the *NOT* operator is applied to the third *qubit* $|\sigma\rangle$ when the current states of the first two *qubits* $|\psi\rangle$ and $|\phi\rangle$ are both $|1\rangle$:

Similarly to QTs of multiple *qubits* which were obtained by the tensor product performed over unitary transformations, Eq.(11) presents the matrix structure of such QT, when $|X\rangle$ is the initial state:

In order to obtain information from a quantum system, it is necessary to apply measurement operators, defined by a set of linear operators M_m , called projections. The index *m* refers to the possible measurement results. If the state of a 1-dimensional quantum system is $|\Psi\rangle$ immediately before the measurement, the probability of an outcome occurrence is given by $p(|\Psi\rangle) = \frac{M_m |\Psi\rangle}{\sqrt{\langle \Psi | M_m^{\dagger} M_m |\Psi \rangle}}$. When measuring

a qubit $|\psi\rangle$ with $\alpha, \beta \neq 0$, the probability of observing $|0\rangle$ and $|1\rangle$ are, respectively, given by the following expressions:

$$p(0)=\langle \phi | M_0^{\dagger} M_0 | \phi
angle = \langle \phi | M_0 | \phi
angle = | lpha |^2;$$

$$p(1) = \langle \phi | M_1^{\mathsf{T}} M_1 | \phi \rangle = \langle \phi | M_1 | \phi \rangle = |\beta|^2$$

After the measuring process, the quantum state $|\psi\rangle$ has $|\alpha|^2$ as the probability to be in the state $|0\rangle$ and $|\beta|^2$ as the probability to be in the state $|1\rangle$.

In multidimensional systems, the operators M_m^n and $p_N(m)$ denote the *m*-projection and corresponding probability measure, both performed on the *n*-qubit.

4 INTERPRETING XOR OPERATORS BASED ON QC

The description of intuitionistic fuzzy sets from the QC viewpoint extends the work in (Mannucci, 2006)

by modeling an element $\tilde{x} = (x_1, x_2)$ by a pair of onedimensional qubit quantum states $(|x_1\rangle, |x_2\rangle)$ where:

$$|x_1\rangle = \sqrt{1-x_1}|0\rangle + \sqrt{x_1}|1\rangle; \qquad (12)$$

$$|x_2\rangle = \sqrt{1-x_2}|0\rangle + \sqrt{x_2}|1\rangle.$$
(13)

By modeling fuzzy operators in QC(Avila et al., 2015), the Product t-norm T_P and Algebriac Sum tconorm can be represented through the *Toffoli* gate (T) and the standard negation through the *Pauli-X* gate (N).

So the first step to generate the quantum representation for both Xor operators \oplus_{I_P} and \otimes_{I_P} is to apply *De Morgan's law* related to t-(co)norms T(S) and fuzzy negation N as present in Eqs. (5) and (6), resulting on the following expressions for their membership and non-membership degrees when considering (\tilde{x}, \tilde{y}) as input:

$$u_{\oplus_{I}} = N(T(N(T(x_{2}, y_{1})), N(T(x_{1}, y_{2}))))$$
(14)
$$T(N(T(X_{1}, y_{2})))$$
(15)

$$v_{\oplus_I} = T(N(T(N(x_1), N(y_2))), N(T(N(x_2), N(y_1))))$$
(15)

$$\mu_{\otimes_{I}} = T(N(T(N(x_{1}), N(y_{1}))), N(T(N(x_{2}), N(y_{2})))) \quad (16)$$

$$\nu_{\otimes_{I}} = N(T(N(T(x_{2}, y_{2})), N(T(x_{1}, y_{1})))) \quad (17)$$

Taking $\tilde{x} = (x_1, x_2)$, $\tilde{y} = (y_1, y_2)$, the initial quantum stated $|\phi\rangle$ is the 10-dimensional quantum register:

$$|\phi\rangle = |x_1\rangle \otimes |x_2\rangle \otimes |y_1\rangle \otimes |y_2\rangle \otimes |0\rangle \otimes |0\rangle \otimes |0\rangle \otimes |0\rangle.$$

Then, the quantum representation of the Xor operator \oplus_{I_p} can be obtained by translating Eqs. (14) and (15), respectively resulting in the following QT compositions:

$$\mu_{\oplus_{i}} = M_{1}^{9} \circ N_{5,6,9} \circ T_{9}^{5,6} \circ N_{5,6} \circ T_{6}^{1,4} \circ T_{5}^{2,3}$$
(18)

$$\nu_{\oplus} = M_1^{10} \circ T_{10}^{7,8} \circ N_{2,3,8} \circ T_8^{2,3} \circ N_{2,3} \circ N_{1,4,7} \circ T_7^{1,4} \circ N_{1,4}$$
(19)

Analogously, considering the initial quantum state $|\phi\rangle$, the quantum representation of the Xor operator \otimes_{I_P} can be obtained by translating Eqs. (16) and (17), respectively resulting in QT compositions given as follows:

$$\mu_{\otimes l} = M_1^9 \circ T_9^{5,6} \circ N_{2,4,6} \circ T_6^{2,4} \circ N_{2,4} \circ N_{1,3,5} \circ T_5^{1,3} \circ N_{1,3}$$
(20)

$$\mathbf{v}_{\otimes \overrightarrow{=}} M_1^{10} \circ N_{10,7,8} \circ \mathbf{T}_{10}^{7,8} \circ N_{7,8} \circ \mathbf{T}_8^{1,3} \circ \mathbf{T}_7^{2,4}$$
(21)

For both quantum representations of membership functions, it is used 10 qubits: 2 pairs for the inputs (\tilde{x}, \tilde{y}) , 4 ancillaries qubits to store intermediate results and 1 pair for the final result. The membership degree obtained is stored on qubit 9 and the nonmembership degree on qubit 10.



4.1 Modelling \oplus_{I_P} Quantum Operator

See in Fig. 1 the quantum circuit for the \bigoplus_{I_P} Xor operator, resulting from the composition of Eq.(18) (from T1 to T5) and Eq. (19) (from T6 to T12) which is anticipating the measure operations.

Columns in Table 1 show the non-void amplitude evolution for the most relevant points of this quantum circuit, with T0 denoting the initial quantum state and T12 is the final quantum state resulting on $\bigoplus_{I_P} X$ or operator obtained from the application of all composition quantum operators. And, in such columns, the changed qubits are highlighted.

After performing the circuit in Fig. 1, the measure operator M_1^9 is applied, that is, on the 9th qubit and related to $|1\rangle$, it has probability

$$p_{9\oplus_I}(1) = x_1 y_2 + x_2 y_1 - x_1 x_2 y_1 y_2,$$

corresponding to the membership degree of \oplus_{I_P} Xor operator, obtained by the μ_{\oplus_I} expression. Therefore,

$$p_{9\oplus_I}(1) = \mu_{\oplus_{I_P}(\tilde{x},\tilde{y})}.$$

Analogous, the resulting measure operator M_0^9 has the probability expressed as follows

$$p_{9\oplus i}(1)(0) = 1 - x_1y_2 - x_2y_1 + x_1x_2y_1y_2$$

corresponding to the complement of the membership degree of $\bigoplus_{I_P}(\tilde{x}, \tilde{y})$, meaning that

$$p_{9\oplus_I}(1)(0) = 1 - \mu_{\oplus_I}(\tilde{x}, \tilde{y}).$$

Thus, the 9^{th} qubit is given by the following expression:

$$|\mu_{\oplus_I}\rangle = \sqrt{1 - \mu_{\oplus_I}(\tilde{x}, \tilde{y})} |0\rangle + \sqrt{\mu_{\oplus_I}(\tilde{x}, \tilde{y})} |1\rangle.$$

In addition, by applying the measure operator M_1^{10} , that is, on the 10^{th} qubit and related to $|1\rangle$, the result quantum state has probability

$$p_{10\oplus_I}(1) = x_1 x_2 + x_1 y_1 + x_2 y_2 + y_1 y_2 - x_1 x_2 y_1 - x_1 x_2 y_2 - x_1 y_1 y_2 - x_2 y_1 y_2 + x_1 x_2 y_1 y_2,$$

corresponding to the non-membership degree of $\bigoplus_{I_P}(\tilde{x}, \tilde{y})$. Then, it means that

$$p_{10\oplus_I}(1) = \mathsf{v}_{\oplus_I}(\tilde{x}, \tilde{y})$$

non-void amplitudes	<i>T</i> 0	T1	T2	T5	<i>T</i> 8	<i>T</i> 11	T12
$(1-x_1)(1-x_2)(1-y_1)(1-y_2)$	0000000000	0000000000	0000000000	0000000000	0000000000	0000000000	0000000000
$(1-x_1)(1-x_2)(1-y_1)y_2$	0001000000	0001000000	0001000000	0001000000	000100 1 000	0001001000	0001001000
$(1-x_1)(1-x_2)y_1(1-y_2)$	0010000000	0010000000	0010000000	0010000000	0010000000	0010000100	0010000100
$(1-x_1)(1-x_2)y_1y_2$	0011000000	0011000000	0011000000	0011000000	001100 1 000	0011001100	001100110 1
$(1-x_1)x_2(1-y_1)(1-y_2)$	0100000000	0100000000	0100000000	0100000000	0100000000	0100000100	0100000100
$(1-x_1)x_2(1-y_1)y_2$	0101000000	0101000000	0101000000	0101000000	010100 1 000	0101001 1 00	010100110 1
$(1-x_1)x_2y_1(1-y_2)$	0110000000	0110100000	0110100000	0110100010	0110100010	0110100110	0110100110
$(1-x_1)x_2y_1y_2$	0111000000	01111100000	0111100000	01111000 1 0	011110 1 010	0111101 1 10	0111101111
$x_1(1-x_2)(1-y_1)(1-y_2)$	100000000	100000000	100000000	1000000000	1000001000	1000001000	1000001000
$x_1(1-x_2)(1-y_1)y_2$	1001000000	1001000000	1001010000	10010100 1 0	100101 1 010	1001011010	1001011010
$x_1(1-x_2)y_1(1-y_2)$	1010000000	1010000000	1010000000	1010000000	101000 1 000	1010001 1 00	101000110 1
$x_1(1-x_2)y_1y_2$	1011000000	1011000000	10110 1 0000	10110100 1 0	101101 1 010	1011011 1 10	101101111 1
$x_1x_2(1-y_1)(1-y_2)$	1100000000	1100000000	1100000000	1100000000	110000 1 000	1100001 1 00	110000110 1
$x_1x_2(1-y_1)y_2$	1101000000	1101000000	1101010000	11010100 1 0	110101 1 010	11010111110	1101011111 1
$x_1x_2y_1(1-y_2)$	1110000000	1110100000	1110100000	11101000 1 0	111010 1 010	1110101 1 10	1110101111 1
$x_1 x_2 y_1 y_2$	1111000000	11111100000	11111110000	11111100 1 0	111111 1 010	1111111 1 10	1111111111 1

Table 1: Evolution of superposition quantum registers in modelling quantum circuit: $\bigoplus_{I_P}(\tilde{x}, \tilde{y})$.

In analogous way, the resulting measure operator M_0^{10} has the probability

$$p_{10\oplus_{I}}(0) = 1 - x_1 x_2 - x_1 y_1 - x_2 y_2 - y_1 y_2 + x_1 x_2 y_1$$

 $+x_1x_2y_2+x_1y_1y_2+x_2y_1y_2-x_1x_2y_1y_2,$

corresponding to the complement of the nonmembership degree of $\bigoplus_{I_P}(\tilde{x}, \tilde{y})$, that is,

$$p_{10\oplus_I}(0)=1-\mathsf{v}_{\oplus_I}(\tilde{x},\tilde{y}).$$

Thus, the 10^{th} qubit is given as follows:

$$|\mathbf{v}_{\oplus_I}\rangle = \sqrt{1 - \mathbf{v}_{\oplus_I}(\tilde{x}, \tilde{y})} |0\rangle + \sqrt{\mathbf{v}_{\oplus_I}(\tilde{x}, \tilde{y})} |1\rangle.$$

Concluding, the interpretation of intuitionistic fuzzy values related to the \oplus_{I_P} Xor connective is provided by the pair $(|\mu_{\oplus_I}\rangle, |\nu_{\oplus_I}\rangle)$ of quantum registers. **Remark 4.1.** As a relevance, the entanglement of such qubits makes the fuzzy quantum circuits different from circuits modeling other logical approaches.

For instance, by taking the results from $\oplus_{I_P} Xor$ operator one can easily observe that when the measurement applied to the 10^{th} qubit is related to $|1\rangle$, then it returns 7^{th} and 8^{th} qubits also related to $|1\rangle$, meaning that these three qubits collapse to value 1.

This is a phenomenon that does not occur in standard fuzzy set theory: the result of measurement (observation) affects the state of other arguments.

So, these values can be used in next calculations performed by other functions but involving such qubit systems, independently of the reuse of circuit inputs, which are also restored from 1^{rs} to 4^{th} qubits.

4.2 Modelling \otimes_{I_P} Quantum Operator

Fig. 2 shows the \otimes_{I_P} quantum circuit related to composition of Eq.(20) (from T1 to T7) and Eq.(21) (from T8 to T12) without the measure operations.

Analogously, based on the entanglement of such qubits, one can easily observe that when the measurement 10^{th} qubit is related to $|0\rangle$ simultaneously, the 5^{th} and 6^{th} qubits are also related to $|0\rangle$, meaning that these three qubits collapse to value 0.

Moreover, see columns in Table 2 presenting the evolution of the non-void amplitudes for the most relevant points of this circuit, with T0 being the initial quantum state and T12 is the final quantum state, ie, the quantum state resulting after the application of all the quantum operators related to the $\otimes_{I_P}(\tilde{x}, \tilde{y})$.

After executing this circuit, applying the measure operator M_1^9 , that is, on the 9th qubit and related to $|1\rangle$, it has the following distribution of probability:

$$p_{9\otimes_{I_{p}}}(1) = x_{1}x_{2} + x_{1}y_{2} + x_{2}y_{1} + y_{1}y_{2} - x_{1}x_{2}y_{1} - x_{1}x_{2}y_{2} - x_{1}y_{1}y_{2} - x_{2}y_{1}y_{2} + x_{1}x_{2}y_{1}y_{2}.$$

corresponding to the membership degree of $\otimes_{I_P}(\tilde{x}, \tilde{y})$, then we obtain that

$$p_{9\otimes_I}(1) = \mu_{\otimes_{I_P}(\tilde{x},\tilde{y})}.$$
(22)

Analogously, the measure M_0^9 has the following probability

$$p_{9\otimes_{I}}(0) = 1 - x_{1}x_{2} - x_{1}y_{2} - x_{2}y_{1} - y_{1}y_{2} + x_{1}x_{2}y_{1} + x_{1}x_{2}y_{2} + x_{1}y_{1}y_{2} + x_{2}y_{1}y_{2} - x_{1}x_{2}y_{1}y_{2},$$

which means that the complement of the membership degree of \otimes_{I_P} Xor operator is given as follows

$$p_{9\otimes_I}(0) = 1 - \mu_{\otimes_{I_P}}(\tilde{x}, \tilde{y}).$$

$$(23)$$

Thus, By Eqs. (22) and (23) we have that the 9^{th} qubit can be given as follows:

$$|\mu_{\otimes_{I}}\rangle = \sqrt{1 - \mu_{\otimes_{I}}(\tilde{x}, \tilde{y})} |0\rangle + \sqrt{\mu_{\otimes_{I}}(\tilde{x}, \tilde{y})} |1\rangle.$$

non-void amplitudes	<i>T</i> 0	T3	<i>T</i> 6	<i>T</i> 7	<i>T</i> 8	<i>T</i> 11	T12
$(1-x_1)(1-x_2)(1-y_1)(1-y_2)$	0000000000	0000000000	0000000000	0000000000	0000000000	0000000000	0000000000
$(1-x_1)(1-x_2)(1-y_1)(1-y_2)$	0000000000	0000000000	0000000000	0000000000	0000000000	0000000000	0000000000
$(1-x_1)(1-x_2)(1-y_1)y_2$	0001000000	0001000000	0001010000	0001010000	0001010000	0001010000	0001010000
$(1-x_1)(1-x_2)y_1(1-y_2)$	0010000000	0010100000	0010100000	0010100000	0010100000	0010100000	0010100000
$(1-x_1)(1-x_2)y_1y_2$	0011000000	0011100000	00111110000	00111100 1 0	0011110010	0011110010	0011110010
$(1-x_1)x_2(1-y_1)(1-y_2)$	0100000000	010000000	0100010000	0100010000	0100010000	0100010000	0100010000
$(1-x_1)x_2(1-y_1)y_2$	0101000000	0101000000	0101010000	0101010000	010101 1 000	0101011000	010101100 1
$(1-x_1)x_2y_1(1-y_2)$	0110000000	0110100000	01101 1 0000	01101100 1 0	0110110010	0110110010	0110110010
$(1-x_1)x_2y_1y_2$	0111000000	01111100000	01111110000	01111100 1 0	011111 1 010	0111111010	0111111011
$x_1(1-x_2)(1-y_1)(1-y_2)$	100000000	1000100000	1000100000	1000100000	1000100000	1000100000	1000100000
$x_1(1-x_2)(1-y_1)y_2$	1001000000	1001100000	1001110000	1001110010	1001110010	1001110010	1001110010
$x_1(1-x_2)y_1(1-y_2)$	1010000000	1010100000	1010100000	1010100000	1010100000	1010100 1 00	101010010 1
$x_1(1-x_2)y_1y_2$	1011000000	1011100000	10111110000	1011110010	1011110010	1011110110	101111011 1
$x_1x_2(1-y_1)(1-y_2)$	1100000000	1100100000	1100110000	11001100 1 0	1100110010	1100110010	1100110010
$x_1x_2(1-y_1)y_2$	1101000000	1101100000	11011 1 0000	11011100 1 0	110111 1 010	1101111010	110111101 1
$x_1x_2y_1(1-y_2)$	1110000000	1110100000	1110110000	11101100 1 0	1110110010	1110110 1 10	111011011 1
$x_1 x_2 y_1 y_2$	1111000000	11111100000	11111110000	11111100 1 0	111111 1 010	1111111 1 10	1111111111 1

Table 2: Evolution of superposition quantum registers in modelling quantum circuit: $\otimes_{I_p}(\tilde{x}, \tilde{y})$.

Moreover, applying the measure operator M_1^{10} , that is, measurement on the 10^{th} qubit and related to $|1\rangle$, the resulting probability is given as follows:

 $p_{10\otimes I_{P}}(1) = x_{1}y_{1} + x_{2}y_{2} - x_{1}x_{2}y_{1}y_{2},$

which corresponds to the non-membership degree of $\otimes_I(\tilde{x}, \tilde{y})$, therefore, we obtain that

$$p_{10\otimes_I}(1) = \mathbf{v}_{\otimes_I}(\tilde{x}, \tilde{y}). \tag{24}$$

Analogous, the measure M_0^{10} has the probability

$$p_{10\otimes x}(0) = 1 - x_1y_1 - x_2y_2 + x_1x_2y_1y_2,$$

corresponding to the complement of the nonmembership degree of $\otimes_{I_P}(\tilde{x}, \tilde{y})$, that is,

$$p_{10\otimes_I}(0) = 1 - \mathbf{v}_{\otimes_I}(\tilde{x}, \tilde{y}).$$
 (25)

And finally, by Eqs. (24) and (25) we have that the 10^{th} qubit is given by:

$$|\mathbf{v}_{\otimes_{I}}\rangle = \sqrt{1 - \mathbf{v}_{\otimes_{I}}(\tilde{x}, \tilde{y})}|0\rangle + \sqrt{\mathbf{v}_{\otimes_{I}}(\tilde{x}, \tilde{y})}|1\rangle.$$

Concluding, the interpretation of intuitionistic fuzzy values related to the \otimes_{I_P} Xor connective is provided by the pair $(|\mu_{\otimes_I}\rangle, |v_{\otimes_I}\rangle)$ of quantum registers.

5 CONCLUSIONS

This paper describes the interpretation of two classes of Xor operations on A-IFS through concepts of QC. It was modelled using a quantum register using operations over fuzzy sets described by QTs.

Therefore, the presented approach to interpretation of intuitionistic fuzzy valued from quantum registers and quantum states shows another basic construction in the specification of fuzzy expert systems from QC, in order to obtain new information technologies based on intuitionistic fuzzy approach.

Computer systems based on A-IFL and performed over quantum computers may be able to generate an output considering the manipulation of inaccurate data and also dealing with imprecision in the model of rule-based system, by taking advantage of properties as quantum parallelism.

Further work aims at consolidation of this specification including not only other fuzzy connectives but also constructors (e.i. automorphisms and reductions) and the corresponding extension of (de)fuzzyfication methodology from formal structures provided by QC.

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