

The Devil Lies in the Details for Species Coexistence Stability in Ablated and Unablated Five-Species Evolutionary Spatial Cyclic Games

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Abstract: I present exploration of key results from the series of ablated five-species “Rock-Paper-Scissors-Lizard-Spock” minimal agent-based evolutionary models of biodiversity introduced by Zhong, Zhang, Li, Dai, & Yang in their 2022 paper “Species coexistence in spatial cyclic game of five species” (*Chaos, Solitons and Fractals*, 156: 111806). At the heart of Zhong et al.’s model of ecosystems coexistence is the *Elementary Step* (ES) algorithm in which one or two neighboring agents are chosen at random to engage in one or more interactions selected at random from the set {COMPETE, REPRODUCE, MOVE}. Minor revisions to the ES algorithm have recently been introduced to make it more computationally efficient in space and in time, and one contribution of this paper is to demonstrate that switching to this “Revised ES” (RES) has the unexpected effect of totally changing the outcomes of Zhong et al.’s simulation experiments. I present analysis of the RES-based experiments which shows that the key difference is that in RES the likelihood of an agent moving is decoupled from the likelihood of an agent reproducing or competing, whereas in the original ES the likelihoods of the three possible actions are interdependently coupled. The fact that such relatively minor changes to the ES algorithm result in such major changes in the experiment outcomes casts significant doubt on the extent to which results such as those from Zhong et al.’s original experiments can be trusted as truly representative of the real-world biological systems that they are supposedly intended to model. Python source-code available on GitHub can be used to replicate the results presented here.

1 INTRODUCTION


1.1 Evolutionary Spatial Cyclic Games

There is a long tradition of empirical bioinformatics research exploring issues in ecosystems biodiversity, and specifically in the population dynamics of species co-existence, by running sets of computer simulation experiments of minimal models of multiple interacting species in what are known as *Evolutionary Spatial Cyclic Games* (ESCGs). In these simulated systems, each individual agent (i.e., each modelled organism) in the population can *move* (i.e., each agent i occupies a specific position in space at time t , denoted by $\vec{p}_i(t)$, and it may move to some different location later in time such that $\vec{p}_i(t + \Delta_t) \neq \vec{p}_i(t)$). Neighboring agents in these models may *compete* in pairwise fight-to-the-death predator/prey interactions, in which the loser dies leaving an empty space; and any empty

space may subsequently be filled with a new agent, created by *reproduction*, asexually spawning a clone of a neighboring living agent into the empty spot.

Inspired by the work of (May and Leonard, 1975), the complexities of real-world predator-prey interactions are reduced in these models to very simple two-player competitive games that require almost no time to simulate, and to give interesting coevolutionary population dynamics the game is usually *cyclic*, i.e. it has intransitive dominance relationships. The most well-known of such cyclic games is the hand-gesture game Rock-Paper-Scissors (RPS): *Rock* beats *Scissors*; *Scissors* beats *Paper*; and *Paper* beats *Rock*. Using RPS as the predator-prey game allows for the study of three-species model ecosystems (i.e., where there are three species denoted by s_R , s_P , and s_S , with individuals in species s_R always playing *Rock*, etc).

Despite their abstract simplicity, RPS-based ESCGs have been useful in explaining aspects of real-world biological systems in which such cyclic dominance relationships do occur and are sustained:

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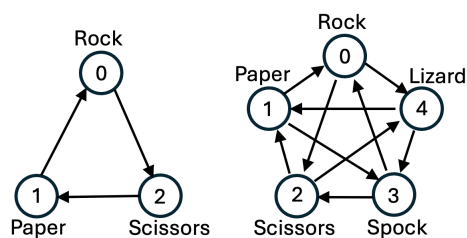


Figure 1: Dominance network digraphs for Rock-Paper-Scissors (RPS: left) and for Rock-Paper-Scissors-Lizard-Spock (RPSLS: right). In both digraphs, the N_S different species S_i are denoted by nodes labelled with index $i \in \{1, \dots, N_S - 1\}$, with arcs (directed edges) running from the dominator (“predator”) species to the dominated (“prey”) species. The rules of RPSLS are: scissors cut paper; paper covers rock; rock blunts scissors; scissors decapitates lizard; lizard eats paper; paper disproves Spock; Spock vaporizes rock; rock crushes lizard; lizard poisons Spock; and Spock smashes scissors.

for examples, see the various papers cited in (Nag Chowdhury et al., 2023).¹ An extension of RPS to a five-gesture cyclic game known as Rock-Paper-Scissors-Lizard-Spock (RPSLS) was introduced by (Kass and Bryla, 1998) and that game has since been widely used to study five-species model ecosystems. Groundbreaking papers in this field of study were published over the period 2006-2008: see, for example, (Laird and Schamp, 2006; Reichenbach et al., 2007a; Reichenbach et al., 2007b; Laird and Schamp, 2008; Reichenbach et al., 2008)

This current paper presents results from ESCG experiments where one or more of the dominance relationships in the RPSLS game have been removed, giving rise to systems that I refer to as ESCGs with *ablated dominance networks*.

1.2 Ablated Dominance Networks

Let N_S denote the number of species in the ESCG model. The intransitive predator/prey dominance relationships of the $N_S=3$ RPS and $N_S=5$ RPSLS games are illustrated as directed graphs (or *digraphs*) in Figure 1. Such digraphs are referred to here as *dominance networks*.

In a recent paper, (Zhong et al., 2022) explored the effects of selectively ablating one or more relationships from the RPSLS dominance network. That is, Zhong et al. deleted one or more of the arcs (directed edges) from the RPSLS digraph, and then observed the frequency distribution of asymptotic experiment outcomes from large numbers of independent and

identically distributed (IDD) runs of the ESCG simulator using the ablated dominance network. Here, “asymptotic” is interpreted in its usual sense, of the response of the system as time $t \rightarrow \infty$, and the experiment outcomes of interest are the number of co-existing species at the end of the experiment. If we let T denote the pre-specified experiment end-time (where T is assumed to be sufficiently big that the system does indeed settle to its asymptotic state), and if we let $n_s(t)$ denote the number of species in the system still surviving at time t , then the key observations in Zhong et al.’s study was the set of frequencies $F(n_s(T)=c) : c \in \{0, 1, \dots, N_S\}$.

Zhong et al. studied the effect of varying the *mobility* parameter in the ESCG (denoted by M , and explained in more detail in Section 2), which determines how mobile the individual agents are, and for each ablated dominance network they plotted graphs of $F(n_s(T)=c)$ as a function of M . Zhong et al. reported a diversity of experiment outcomes, depending on exactly which arc or arcs had been ablated from the RPSLS dominance network. Section 2 summarises the work of Zhong et al. in more detail.

In (Cliff, 2024b) I presented results which replicated and confirmed Zhong et al.’s results, before questioning whether they had chosen a sufficiently large value of T for their simulations to truly converge on asymptotic states. In this paper I go on to fully explore the effects of re-running Zhong et al.’s experiments with a slightly different style of ESCG simulation recently introduced in (Cliff, 2024a), which involves what are *prima facie* only very minor adjustments to the core algorithm in the ESCG simulation, the so-called *Elementary Step* (ES). The initial motivation for altering the ES was to improve its efficiency in space and in time. The revised algorithm and associated initial results were first published in (Cliff, 2024a), a paper presented at an international conference on modeling and simulation methods, where the focus was on the computational efficiency increases (reductions in execution time for the simulations) and where I did not seek to explore the effects of using the revised ES algorithm across the full range of ablated dominance-network species co-existence experiments reported by Zhong et al..

My new results and analysis presented here show that switching from the “Original” Elementary Step (OES) used by (Zhong et al., 2022) to the “Revised” Elementary Step (RES) introduced in (Cliff, 2024a) has a major effect on every single one of the key experiment outcomes reported in (Zhong et al., 2022); that is, for *all* of Zhong et al.’s key results, I demonstrated in (Cliff, 2024b) that they can be replicated very well with OES and then in this paper, when I

¹These include (Jackson and Buss, 1975; Sinervo and Lively, 1996; Durrett and Levin, 1998; Elowitz and Leibler, 2000; Kerr et al., 2002; Gilg et al., 2003; Lankau and Strauss, 2007; Nahum et al., 2011) and (Guill et al., 2011).

switch the simulator to instead use RES, I get totally different outcomes – identifying and explaining that difference is the key contribution of this paper. The Python source-code initially developed for the research reported in (Cliff, 2024a) is freely-available public-domain open-source software on GitHub,²

I give further details of the OES-based ESCG as used by Zhong et al. (and many other researchers) in Section 2, and then explain the differences between OES and RES in Section 3. I discuss differences between OES and RES results in Section 4. Then in Section 5 I explain the root cause of those differences and discuss the extent to which they cast doubts on just how far results from OES can be trusted as being representative of the real-world biosystems that they are intended to be informative models of.

2 THE ESCG ALGORITHM

Zhong et al. used the *de facto* standard algorithm for their Evolutionary Spatial Cyclic Game (ESCG), the core elements of which were established in (Reichenbach et al., 2007a).

The position \vec{p}_i of each individual agent i in the population is represented by the values on a Cartesian lattice l of some dimensionality, and almost all work in this field studies systems where the lattice is a 2-dimensional square of side-length L , so the maximum total population size is $N=L \times L$, where the position \vec{p}_i of an individual agent i can be represented by an (x, y) row-column co-ordinate pair indexing into the lattice as $l(x, y)$. Cells (or ‘positions’ or ‘locations’) in the lattice can be empty, denoted here by \emptyset , or they can be occupied by a single individual agent. In the RPSLS-based systems explored by (Zhong et al., 2022), the number of species is fixed at $N_S=5$, and the species identifier indexes will here run $0, \dots, N_S - 1$, so the set of possible lattice cell-values is $C_S = \{\emptyset, 0, \dots, 4\}$.

Let $\mathcal{U}(S)$ denote the uniform (equiprobable) random choice of a member from set S , and let $\mathcal{U}[n_{lo}, n_{hi}]$ denote a fresh random draw from a uniform distribution over the range $[n_{lo}, n_{hi}] \subset \mathbb{R}$. To initialise the lattice, some researchers include a number of empty cells in the initial state, but (Zhong et al., 2022) do as many other authors do, starting instead with every cell containing an agent, with each species s_i being chosen with probability $\frac{1}{N_S} \forall i$ such that $l(x, y) = \mathcal{U}(\{0, \dots, 4\}) \forall (x, y)$.

The ESCG is then simulated as a simple discrete-time system. At the core of the simulation is the *Elementary Step* (ES), the original and revised forms

of which (i.e., the OES and RES) are discussed further in Section 3. In both the OES and the RES, the ES implements an interaction between two randomly-chosen neighboring cells, and this requires the researcher to commit to a *neighborhood function* and to a specific *boundary condition*. The neighborhood function is typically either the von Neumann or the Moore. The boundary condition is typically either *periodic*, where the square lattice is, topologically, a toroidal surface; or *no-flux*, where the edges of the lattice are hard boundaries on a ‘walled garden’. Zhong et al. used the von Neumann Neighborhood and no-flux boundary conditions, and so that is what I used in the replication experiments described here.

Implementing the random selection of a cell and its randomly chosen neighbor is trivial. Once two neighboring cells have been chosen, they may interact, and the opportunity set of possible interactions is the same for both OES and RES:

- If both cells contain agents, they might *compete* according to the rules of the ESCG’s dominance network for species-to-species interactions. Competition will result either in a draw (no change), or the dominated ‘prey’ individual dies and the lattice cell it was occupying is replaced by \emptyset .
- If one cell contains an agent and the other contains \emptyset (i.e., is empty), the agent might *reproduce* by cloning a copy of itself into the empty cell.
- The contents of the two cells might be swapped, which implements a one-step *move* for each of the two cells’ contents.

Which, if any, of these three interactions occurs on any one ES is determined stochastically, and the random choice of interaction is determined by three system hyperparameters: μ , which primarily affects the probability of choosing *compete*; σ , which primarily affects the probability of choosing *reproduce*; and ϵ which primarily affects the probability of choosing *move*. Zhong et al. describe these three as “probabilities” but (as is explained below in Section 3) this is incorrect: they are relative *interaction rates*.

The convention in this field of research is to talk of time-steps at two levels of granularity: the ES is the finer-grained time-step, but it involves only a single potential interaction between two randomly chosen cells in what is usually a large lattice of such cells. At a coarser-grained level, researchers talk of the *Monte Carlo Step* (MCS) as being the primary unit of time, a time-step in which all N cells in the lattice can potentially experience at least one interaction: i.e., one MCS involves N elementary steps being executed.

In all the experiments they reported in detail, Zhong et al. used $L=200$ and (in common with many

²See github.com/davecliff/ESCG_Python

other researchers) they set $\mu=\sigma=1.0$ for all experiments, stating that they did so without loss of generality. Following the *de facto* standard procedure established in many prior papers in this field, Zhong et al. did not explicitly set values for the movement factor ϵ , but instead systematically varied, as the primary independent variable in their experiments, a *mobility* parameter M , which relates the movement rate ϵ to the lattice size N via $\epsilon=2MN$. For each of their key experiments, Zhong et al. varied M in discrete steps over $M \in [10^{-7}, 10^{-3}]$, sampling the response of the system, its frequency distribution of final outcomes, at each M -value. At each value of M sampled, Zhong et al. performed $N_{\text{IID}}=500$ independent and identically distributed (IID) simulation experiments.

Finally, to complete the specification for one of the RPSLS-style ESCGs simulated by (Zhong et al., 2022), the dominance network needs to be defined. In (Cliff, 2024b) I introduced terminology to distinguish the various ablated RPSLS dominance networks that Zhong et al. explored (each of which, Zhong et al. referred to solely by reference to which figure-number in their paper the relevant network diagram appeared in): let $Z0$ denote the original unablated RPSLS network as illustrated in Figure 1, and then let any version of that network with n_a arcs ablated be referred to as a Zn_a network, so for example $Z0$ with two ablated arcs would be referred to as a $Z2$ network; then in the cases where Zhong et al. studied two topologically nonisomorphic networks of type Zn_a , add qualifying letters a and b as suffixes. In total, Zhong et al. studied the dynamics of ESCGs with six different ablated versions of $Z0$, referred to as $Z1$, $Z2a$, $Z2b$, $Z3a$, $Z3b$, and $Z4$; for full details, see (Cliff, 2024b).

The ESCG algorithm used by Zhong et al. is simple to describe: the 2D lattice L is randomly initialised at $t=0$, as described above, and the dominance network \mathcal{D} is set up. An outer time-step loop iterates over each MCS, within which an inner loop iterates N times over the Elementary Step (ES) procedure ELSTEP. After each MCS time-step, the population density of each of the N_S species (denoted by ρ_i for $i \in \{0, \dots, N_S - 1\}$), along with the density of empty cells (denoted by ρ_0) is recorded, optionally along with the state of the lattice at that MCS. At the end of each IID simulation run, the time-series of densities ρ_0 and $\rho_i \forall i$ are written to an output file, and the set of frequency distributions of experiment outcomes (i.e., $F(n_s(T)=c)$ for $c \in \{1, \dots, 5\}$) are updated. The final F s for the set of N_{IID} simulations is, in the current context, the primary data of interested returned by the ESCG.

The elementary step (ES) algorithm ELSTEP is similarly straightforward. It is trivial to randomly

choose a cell in the lattice, and to randomly choose a neighbor of that cell given the pre-defined neighborhood function and boundary conditions. Implementing the three individual cell-to-cell interaction procedures COMPETE, REPRODUCE, and MOVE is similarly unproblematic. However, to better facilitate the comparison between OES and RES, there is some subtlety required in describing the method for randomly selecting the interaction(s) to be enacted on any one elementary step: here we'll assume that ELSTEP can call a function GET3PR which returns three *interaction probabilities* P_μ^* , P_σ^* , and P_ϵ^* that respectively set the likelihood of ELSTEP calling COMPETE, REPRODUCE, and MOVE. This can be simply implemented by generating three random real values r_1 to r_3 each via $\mathcal{U}[0, 1]$, and then executing COMPETE if $r_1 < P_\mu^*$; executing REPRODUCE if $r_2 < P_\sigma^*$; and executing MOVE if $r_3 < P_\epsilon^*$. For OES, GET3PR always returns two of the three probabilities as zero, and the other probability as one – this ensures that only one of the three possible interactions occurs on any one elementary step, which is the convention introduced in the seminal papers of (Reichenbach et al., 2007a); a convention that has been continued in all published literature that replicates and builds upon those works.

The calculation of the three interaction probabilities is where the revised elementary step (RES) differs from the original elementary step (OES), and that issue is discussed in more detail in the next section.

3 OES vs. RES

The difference between the OES, used by (Zhong et al., 2022) and many others,³ and the RES, introduced in (Cliff, 2024a), boils down to how the three probabilities used in ELSTEP, P_μ^* , P_σ^* , and P_ϵ^* , are calculated from the interaction rates μ , σ , and ϵ . Because we're repeatedly dealing here with triplets of values, let $\vec{R} = (\mu, \sigma, \epsilon)$ denote the vector of interaction rates; let $\vec{P}^*(t_{ES}) = (P_\mu^*(t_{ES}), P_\sigma^*(t_{ES}), P_\epsilon^*(t_{ES}))$ denote the vector of interaction probabilities at time t_{ES} , measured here at the micro-level of the count of individual elementary steps, rather than the macro level of MCS-number; and let $\vec{P} = (P_\mu, P_\sigma, P_\epsilon)$ denote

³The core of OES can be traced back to (Gillespie, 1976), cited approvingly in the landmark paper of (Reichenbach et al., 2007a). Other authors who have reported results from OES-based studies include (Nagatani et al., 2018; Kabir and Tanimoto, 2021; Mood and Park, 2021; Park, 2021; Avelino et al., 2022; Bazeia et al., 2022; Menezes et al., 2022a; Menezes et al., 2022b; Zhang et al., 2022; Park and Jang, 2023; Menezes et al., 2023); and (Kubyana et al., 2024).

the vector of normalized probabilities, explained further below.

The OES method, which is the *de facto* standard method in ESCG modelling, can be traced back to first usage in the landmark papers of (Laird and Schamp, 2006) and (Reichenbach et al., 2007a). In those original sources, each of the elements of \vec{R} clearly have the potential to take on values greater than one.

However, quite often the three interaction rates have since been described by other authors of papers using OES as being bounded to the range $[0.0, 1.0]$ and are also referred to by several authors as ‘probabilities’, despite the fact that they are not actually treated as literal probabilities: for example, $\mu = 0.5$ does not necessarily mean there is a 50% chance of COMPETE being the next interaction: it might do, and it might not, depending on the contemporaneous values of σ , and ϵ . This is because the three interaction rates are first subjected to a *normalization* process that divides each of the three by the sum of the three, giving three “normalized probabilities” P_μ, P_σ , and P_ϵ , the elements of \vec{P} , which are each within $[0, 1]$ and which collectively sum to one.

To understand this, the values for the interaction rates μ , σ , and ϵ can be thought of in OES as indications for what relative proportion of a three-segment roulette-wheel is allocated to each possible interaction, and then in the OES the conceptual roulette-wheel is spun once, generating a random value $r(t_{ES}) = \mathcal{U}[0.0, 1.0]$, which is then compared to the elements of \vec{P} such that the interaction at t_{ES} will be COMPETE if $0 \leq r(t_{ES}) < P_\mu$; REPRODUCE if $P_\mu \leq r(t_{ES}) < P_\mu + P_\sigma$; and MOVE otherwise (because $P_\epsilon = 1.0 - (P_\mu + P_\sigma)$). This method is widely known as *roulette selection* in the literature on evolutionary computation (see e.g. (Lloyd and Amos, 2017)). Hence, with OES, $\vec{P}^*(t_{ES})$ is always either $(0, 0, 1)$ or $(0, 1, 0)$ or $(1, 0, 0)$, roulette-selected on the basis of \vec{P} and $r(t_{ES})$: i.e., only one randomly-chosen type of interaction can ever happen per single OES.

The corresponding calculation in RES is much simpler to explain because the interaction rates μ and σ are interpreted instead as actual probabilities, defined only over $[0, 1]$, and mapping directly (via the identity function) onto P_μ^* and P_σ^* , respectively; the movement rate ϵ is required to be non-negative and is interpreted as mapping directly onto P_ϵ^* , with the simple filter that if $\epsilon > 1.0$ then it is capped to one. Thus in RES there is no \vec{P} and instead an essentially direct mapping from \vec{R} to $\vec{P}^*(t_{ES})$. This means that, whereas in OES there is a constraint that $|\vec{P}^*(t_{ES})| = 1 \forall t_{ES}$, in RES instead $|\vec{P}^*(t_{ES})| \in [0, 3] \in \mathbb{R}$ – i.e., at the limit, $\vec{P}^*(t_{ES}) = (1, 1, 1)$ is perfectly well defined.

On the face of it, the change between OES and

RES is pretty minimal. Whereas OES’s roulette-selection means that only one interaction can occur per ES, the definition of RES allows up to three interactions to occur per ES. For example, if $\mu = \sigma = \epsilon = 1.0$ (i.e., $\vec{R} = (1, 1, 1)$) then in OES one of the three possible actions will be randomly chosen, using $r(t_{ES})$ and $\vec{P} = (\frac{1}{3}, \frac{1}{3}, \frac{1}{3})$, and then $\vec{P}^*(t_{ES}) \in \{(1, 0, 0), (0, 1, 0), (0, 0, 1)\}$ depending on $r(t_{ES})$; whereas in RES, for the same \vec{R} we have $\vec{P}^*(t_{ES}) = \vec{R} = (1, 1, 1)$ and hence, whenever possible, on each RES there will be a call to COMPETE, followed by a call to REPRODUCE, followed by a call to MOVE. This means that, when $\vec{R} = (1, 1, 1)$, if possible a single agent can, in one RES, kill its neighboring opponent in COMPETE, and then REPRODUCE a clone of itself into the \emptyset -cell just vacated by the dead opponent; and then MOVE into the neighboring cell, swapping position with its newly-spawned offspring. Whereas with OES the same agent would engage in at most one randomly-chosen interaction per ES, and so would require three consecutive ESs to achieve the same sequence of events.

For that reason, one might expect a RES-based replication of any of the experiments reported in (Zhong et al., 2022) to proceed up to three times faster than the comparable OES-based experiment – and indeed it was a desire for reduced run-times that motivated the development of RES, as reported in (Cliff, 2024a), where further analysis and discussion of the speed-up enabled by RES is given. Here, the focus is not on the run-times of the simulations, but instead on the unexpected results that I report in the next section: the outcomes of every single one of Zhong et al.’s ablation experiments is dramatically altered when the ES is switched from OES to RES, an unexpected consequence of RES that was not explored in any depth in (Cliff, 2024a).

4 RESULTS

4.1 Original Elementary Step (OES)

As with the original results of (Zhong et al., 2022) and also with my replication of those results reported in (Cliff, 2024a), in all the experiments I discuss in (Cliff, 2024b) I set $T = 10^5$ MCS, $L = 200$, $\mu = \sigma = 1.0$, and ranged M from $M_{\min} = 10^{-7}$ to $M_{\max} = 10^{-3}$, using a multiplier of $10^{1/5}$ to give five sampled M data-points per order of magnitude, and I ran OES-based simulation experiments for all six of Zhong et al.’s networks Z1, Z2a, Z2b, Z3a, Z3b, and Z4 – the complete set. All my results presented in (Cliff, 2024b)

agree very well with the original results of (Zhong et al., 2022).

In each set of experiment outcomes the key observable feature, as highlighted by Zhong et al., is a sudden transition part-way between $M=10^{-5}$ and $M=10^{-4}$ from three-or-more-species coexistence being the dominant outcome, to two or fewer coexisting species always being the outcome – i.e., a “collapse” in ecosystems diversity. If this result is general, i.e. if it is applicable to real-world ecosystems and is not just an artefact of the particular details of Zhong et al.’s ESCG modelling methods, then surely we should expect to see it also when the only change made to the ESCG is to switch from OES to RES – and in fact, if anything, we should expect to see the collapse sooner (after a shorter T) when working with RES, because of its potential for a three-times speed-up.

4.2 Revised Elementary Step (RES)

For my RES experiments, I swept M within the range $[5 \times 10^{-8}, 1.25 \times 10^{-5}]$ to give the same range of P_ϵ values as were used by (Zhong et al., 2022) in their OES experiments.

To avoid what would otherwise surely be a strong contender for the most boring set of graphs ever published in a scientific paper, here I give only verbal descriptions of the outcomes of my RES-based replications of Zhong et al.’s experiments. I only need to describe two outcomes because my results were unambiguously bimodal: for networks Z0, Z2b, Z3b, and Z4, the F vs M outcome was essentially identical: the system never deviated from $F(n_s(T)=5)=1.0$ across the entire range of M values; and for the remaining networks (i.e., Z1, Z2a, and Z3a) the outcome was again essentially identical, but in these three cases the system’s asymptotic state was $F(n_s(T)=3)=1.0$ across the entire range of M values. That is, in all my RES experiments, the outcome at T was always either three-species coexistence or five-species coexistence, for every network studied by Zhong et al.: that is, once you switch to RES, the collapse in ecosystems diversity seen every time when using OES simply disappears.

And yet that switch, from roulette-selecting a single interaction in OES to probabilistically selecting one, two, or three interactions in RES, seems only a very small change, a mere implementation detail intended to speed things up a bit – the set of possible interactions remains unchanged, so why does RES result in such different outcomes to OES? And, unless it can be convincingly argued that OES is somehow a better representation, a better abstraction, of the natural world than is RES, then surely this fragility,

these major changes in outcomes resulting from minor changes to the internal mechanisms, casts significant doubt on the generality and applicability of the OES-based results published by (Zhong et al., 2022), and indeed of the results in any other published papers (and there are many) that show results from OES-based ESCGs. This discrepancy in the results is discussed in more detail in the following section.

5 DISCUSSION

To understand why the RES outcomes differ so markedly from the OES ones, it is necessary to go back and consider the step in OES where the three system hyperparameters, the interaction rates μ, σ , and ϵ , are ‘normalized’ to give internal probabilities P_μ, P_σ , and P_ϵ prior to the roulette-selection of a single action for that ES. The normalization process intertwines the three rates, and they cease to be independent, and yet this seems to be rarely if ever acknowledged in the literature.

For example, sweeping M across a range such as the $[10^{-7}, 10^{-3}]$ used by Zhong et al., is described by Zhong et al. (and by many other authors who also present results from parameter-sweeps on M , this being a common way of exploring ESCGs) as systematically varying *mobility*, i.e. treating M as the independent parameter in the experiments. While altering M definitely does alter the likelihood of movement for the agents in the ESCG, in OES it *also* alters the likelihood of competition, and of reproduction. Specifically, in OES, increasing mobility M increases the probability of movement P_ϵ , but it can also *decrease* the probabilities of competition (P_μ) and reproduction (P_σ). Figure 2 shows the changes in values of the three probabilities over the range of M values used by (Zhong et al., 2022), all at $L=200$, for various values of μ and σ including the $\mu=\sigma=1.0$ case used by (Zhong et al., 2022).

As can be seen from the upper left graph in Figure 2, the $\mu=\sigma=1.0$ case, as M is increased P_ϵ does indeed rise as intended, but as it rises so P_μ and P_σ fall, ending up very close to zero once $M \approx 10^{-3}$. Clearly increasing M not only increases mobility, but also greatly reduces the frequency of competition and of reproduction: at $M=10^{-3}, P_\epsilon=0.976$ while $P_\mu=P_\sigma=0.012$; so movement is roughly 80 times more likely than either reproduction or competition. In contrast, when using RES, increasing M does indeed increase P_ϵ , but it has no effect on P_μ or P_σ because the interaction probabilities are independent.

To illustrate this key difference between OES and RES, Figure 3 shows the unit cube in \mathbb{R}^3 formed by

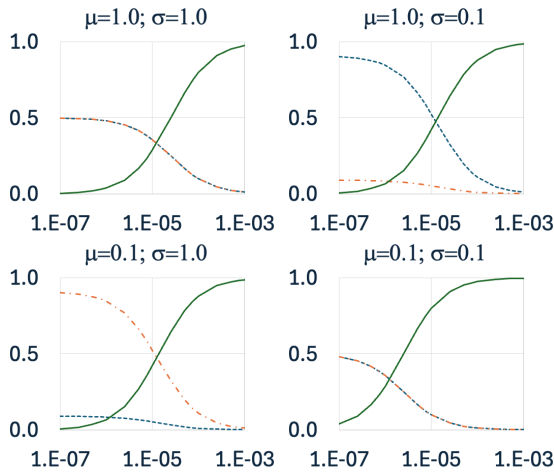


Figure 2: Graphs of the normalized internal probabilities P_μ , P_σ , and P_ϵ for $L=200$ as calculated in the OES for $M \in [10^{-7}, 10^{-3}]$, the range used by (Zhong et al., 2022), for the values of μ and σ shown in each graph title. Each graph shows plots for P_μ (dotted line), P_σ (dash-dot line), and P_ϵ (solid line). In all four graphs the horizontal axis is M while the vertical axis is probability. The upper left graph shows results from $\mu=\sigma=1.0$; upper right shows results from $\mu=1.0, \sigma=0.1$; lower left shows results from $\mu=0.1, \sigma=1.0$; and lower right shows results from $\mu=\sigma=0.1$.

$(P_\mu, P_\sigma, P_\epsilon)$ -space – a cube that I’ll refer to as “ P -space” hereafter. In Zhong et al.’s OES-based experiments with $\mu=\sigma=1.0$. As M is swept from 10^{-7} to 10^{-3} , the three interaction probabilities shift from their initial point close to $(\frac{1}{2}, \frac{1}{2}, 0)$, and move with nonlinearly varying speed along a straight line in P -space, to their final point at approximately $(0, 0, 1)$ – this is represented by the solid blue arrowed line in Figure 3. (At $L=200$, as used by Zhong et al., the actual start-point is $(0.498, 0.498, 0.004)$ and the actual end-point is $(0.01, 0.01, 0.98)$).

However in RES, sweeping M over the corresponding range of values, $P_\mu=P_\sigma=1.0; \forall M$, and the only probability varied as M increases is P_ϵ which rises along the straight line in P -space from its initial point close to $(1, 1, 0)$ to its final point close to $(1, 1, 1)$, which is illustrated by the vertical dashed red arrowed line in Figure 3.

Because the three interaction probabilities are independent in RES, one can readily set P_μ and P_σ to be any pair of values in $[0, 1]^2 \subset \mathbb{R}^2$ and then independently sweep the mobility M across some range which maps P_ϵ to varying over $[0, 1] \subset \mathbb{R}$: that is, in RES it is trivially easy to set up an experiment where the initial point in P -space is anywhere on the floor-plane of the P -space unit cube, and for P_ϵ to then range from zero to one while P_μ and P_σ are held constant, giving a vertical line in P -space. And we need not stop there: in principle, experiments can be run with RES where

across the duration of the experiment the point in P -space moves along some arbitrarily-oriented straight line or curve – there is nothing in RES that limits P_ϵ to being the sole independent variable, nothing that requires only vertical lines in P -space.

However with OES the experimenter’s options are significantly more constrained. To illustrate why this is so, refer back to Figure 2: the upper left graph shows the three interaction probabilities versus M for the case where $\mu=\sigma=1.0$, and the lower right graph shows the same three probabilities for $\mu=\sigma=0.1$; despite the factor of ten difference in the values of μ and σ between these two graphs, the traces shown on the graphs are qualitatively very similar: in both cases, P_μ and P_σ both start at ≈ 0.5 and then decay sigmoidally to ≈ 0.0 , while P_ϵ rises sigmoidally from ≈ 0.0 to ≈ 1.0 .

To understand this, consider first the location of the start-point in P -space for an OES-based experiment where M is swept from some very small initial value (like 10^{-7}) to some larger final value (such as

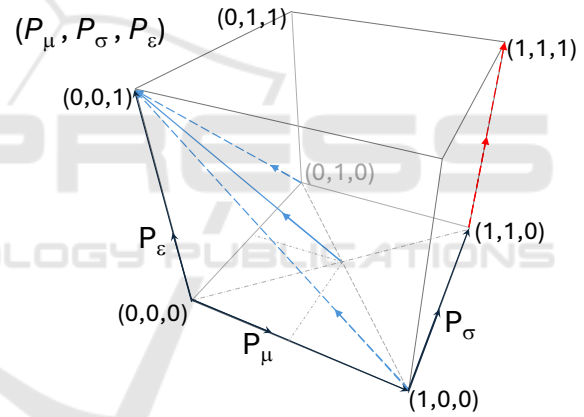


Figure 3: Perspective plot of $(P_\mu, P_\sigma, P_\epsilon)$ -space, referred to here as “ P -space”, which is the unit cube $[0, 1]^3 \subset \mathbb{R}^3$. The floor-plane of the cube is the $(P_\mu, P_\sigma, 0)$ unit square with, from this viewpoint, the P_μ axis running from the origin to lower-right, and the P_σ axis running from the origin to upper right; the vertical axis is P_ϵ . For clarity, seven of the eight corner vertices are labelled with their coordinates; the one unlabelled vertex, nearest the viewpoint, is $(1, 0, 1)$. The solid blue arrowed line from $(0.5, 0.5, 0)$ to $(0, 0, 1)$ shows the straight-line path of P_ϵ values explored by Zhong et al. in their OES-based experiments as they swept M over the range $[10^{-7}, 10^{-3}]$ while keeping $\mu=\sigma=1.0$. The dashed red arrowed line from $(1, 1, 0)$ to $(1, 1, 1)$ shows the straight-line path of P_ϵ values explored when using RES, with M varied over the corresponding range. The dashed blue arrowed lines starting at $(1, 0, 0)$ and at $(0, 1, 0)$, both ending at $(0, 0, 1)$, mark the two vertical-plane edges of the OES unit 3-simplex: its horizontal edge is the line $P_\sigma = 1 - P_\mu$, joining $(1, 0, 0)$ to $(0, 1, 0)$: the significance of the OES unit 3-simplex is discussed further in the text.

10^{-3}). From the way that the three interaction probabilities are calculated as functions of the three interaction rates in the “normalization” step, whenever $\mu=\sigma$, we have $P_\mu=P_\sigma=\frac{\mu}{2\mu+\varepsilon}$. Given that $\varepsilon=2MN$, when $M \ll \frac{1}{2N}$ we have $\varepsilon \approx 0$ and hence $P_\mu=P_\sigma \approx \frac{\mu}{2\mu} = \frac{1}{2}$, for any case when $\mu=\sigma$. Thus it is intuitively obvious that with OES the start-point on the $(P_\mu, P_\sigma, 0)$ ground-plane of the P -space unit cube is determined by the ratio of μ to σ : when $\mu=\sigma=c$ the ratio is 1.0 and the start-point is very close to $(0.5, 0.5, 0.0)$, regardless of the actual value of c ; then, as $\sigma/\mu \rightarrow 0$, the start-point moves ever closer to $(1, 0, 0)$; and as $\sigma/\mu \rightarrow \infty$ the start-point moves ever closer to $(0, 1, 0)$. And, in the two limit cases, when $\mu=0$ we have $P_\mu=0$ but $P_\sigma=\sigma/(\sigma+\varepsilon) \rightarrow 1.0$ as $\varepsilon \rightarrow 0$, giving a start-point of $\approx(0, 1, 0)$; and (*mutatis mutandis*) when $\sigma=0$ so $P_\mu \rightarrow 1.0$ as $\varepsilon \rightarrow 0$, giving a start-point of $\approx(1, 0, 0)$. The key thing to note here is that, for any plausible values of μ and σ , and sufficiently low initial value of M , the start point is *always* on (or very close to) the line $P_\sigma=1-P_\mu$, joining $(1, 0, 0)$ to $(0, 1, 0)$.

Next consider the OES end-point in P -space: this is easy, it’s always going to be at (or very close to) $(0, 0, 1)$, for any plausible values of μ and σ , and sufficiently high final value of M (i.e., for $M \gg \frac{1}{2N}$).

And, taken together, these constraints on the P -space start-points and end-points for OES-based M -sweep experiments mean that the set of possible points sampled in P -space is tightly constrained: the OES-based experiments are limited to the set of points on (or very near) the unit 3-simplex (i.e., the planar triangle defined by the points $(1, 0, 0)$, $(0, 1, 0)$, and $(0, 0, 1)$); thus, OES experiments are simply not set up to explore the response of the system in any area of P -space outside of that three-sided planar area.

To emphasize the importance of this point, let us consider the expected qualitative asymptotic experiment outcomes at each of the seven corner-vertices of P -space for which at least one of the interaction probabilities is nonzero, assuming that the ESCG is run for its duration with the three interaction probabilities held at the vertex:

- **(0, 0, 1):** the agents can move, but they never compete or reproduce, so the whole population of agents wander aimlessly around the lattice, never interacting with one another. This is a dynamic equilibrium, but population-level statistics such as species densities will remain static. There is no chance of species extinctions with the system at this vertex.

- **(0, 1, 0):** agents can reproduce, but not move or compete: any initially empty cells might be filled by reproduction, but once there are no empty cells remaining the system will have converged to a stasis. there’s no chance of extinctions at this vertex.

- **(1, 0, 0):** agents compete, but can’t move or reproduce: any immediate neighbors in the initial lattice configuration who are predators/prey will eventually compete, resulting in the death of the prey. Once all possible predator/prey interactions in the initial lattice configuration have been resolved by competition, the number of empty cells will have increased, but the system will thereafter be at stasis. Species extinctions can occur when the system is at this vertex: the likelihood of extinctions is entirely dependent on the initial distribution of agents on the lattice at $t=0$. With uniform random allocation of agents to the lattice positions at $t=0$ the likelihood of one or more entire species going extinct when the interaction probabilities are held at this vertex is technically nonzero but will be vanishingly small for any sufficiently large value of L to be of serious interest.

- **(0, 1, 1):** agents reproduce and move, but don’t compete: empty cells in the lattice at $t=0$ will eventually be filled by reproduction, and all agents can move around freely; once the last empty cell is filled, the system is thereafter at the same dynamic equilibrium with static species densities as at vertex $(0, 0, 1)$, and so there is no chance of extinctions here either.

- **(1, 0, 1):** agents compete and move, but don’t reproduce: the number of empty cells is likely to rise over the course of the experiment as prey agents are dominated by predators. Species extinctions are entirely likely when the system is at this vertex, the likelihood depending on the dominance network.

- **(1, 1, 0):** agents compete and reproduce, but don’t move: the number of empty cells is likely to rise and fall, with the asymptotic outcome being zero empty cells. Species extinctions can occur at this vertex, the likelihood depending on $l(0)$ (the state of the lattice l at $t=0$) and on the dominance network.

- **(1, 1, 1):** the whole enchilada – agents can compete, reproduce, and move. Species extinctions can occur when the system is at this vertex, the likelihood depending on $l(0)$ and the dominance network.

From this it is clear that, if the ESCG is being run as a minimal model of ecosystems biodiversity and species co-existence, the final three vertices $(1, 0, 1)$, $(1, 1, 0)$, and $(1, 1, 1)$ are of primary interest, because they are the three for which there is a reasonable chance of species extinctions. These three vertices are all on the top-plane of the P -space unit cube, and because none of them are within the unit 3-simplex, they simply can’t be explored via OES-based models. Instead, as mobility M is increased in OES-based models, the system is taken ever-closer to the “aimless wandering” P -space vertex $(0, 0, 1)$.

For example, a naive experimenter might start an OES M -sweep simulation at $\mu=1$ and $\sigma=0$, i.e. at

the (1,0,0) vertex in P -space, and then sweep M over $[10^{-7}, 10^{-3}]$ in the (mistaken) expectation that this will increase only P_e and eventually take the system to the (1,0,1) P -space vertex. However the OES normalization process will instead take the system on a straight-line path toward the (0,0,1) vertex. Some species extinctions may occur en route, and – as I have shown in my replication of the results of (Zhong et al., 2022) — the system will close in on (0,0,1) but P_μ and P_σ can remain nonzero, so extinctions do remain a possibility. But the OES experimenter’s results will say nothing about the asymptotic behavior of the ESCG system at (1,0,1) in P -space: for that, they would need to abandon OES and switch to RES.

The results from RES-based simulations presented here show that when areas of P -space that are unreachable in OES-based simulations are instead explored via the ESCG with RES, the frequency distributions of asymptotic experiment outcomes can change significantly. And surely that then calls into question the generality and the applicability of the results from OES-based simulations to the real-world systems that they are intended as models of: my RES results present a first counter-example to the observation in OES-based systems of the sudden transition in experiment outcomes from being $n_s(T) \geq 3$ at low M to $n_s(T) \leq 2$ at higher M values: this happens as the OES-based system converges on P -space vertex (1,0,0) but does not happen when the RES-based system is instead set to converge on (1,1,1).

And from that it seems reasonable to propose that Zhong et al.’s OES-based results, repeatedly showing collapses in ecosystems coexistence/diversity as mobility increases, are mere artefacts of their decision to use OES in their ESCG model: if they’d used RES, they would have had no story to tell. The only possible line of counterargument that I can think of here would be to defend use of OES on the basis that it is somehow more faithful to ecological reality, more biologically plausible, than RES, and hence to argue that OES should continue to be used in preference to RES because it is, in that sense, the better modelling method. Now maybe I’m biased in favour of my own creation, but I’ve tried very hard to think of how OES could possibly be defended as somehow better than RES, and thus far I have drawn a blank. The two methods strike me as equally abstract, equally (im)plausible approximations to reality, but in OES the experimenter’s hands are tied in P -space, while in RES the experimenter can roam freely within the confines of that cube, which surely opens up more avenues for exploration than repeatedly moving from the base to the apex of the unit simplex, which is seemingly the only option with OES. And, in that

sense, arguably RES is a better approximation to reality than OES, because in many real ecosystems it seems intuitively unlikely that every species’ motility, fecundity, and likelihood of eating or being eaten are interdependently bound together by one function (as in OES), and more likely that these three aspects of species’ existence are to some extent independently variable (as in RES). I am of course very open to conversation with anyone who wants to argue the point: discussion and debate is very welcome here.

In this paper I’ve shown results from RES-based M -sweep experiments where only one path in P -space, along the line from (1,1,0) to (1,1,1), is explored. An obvious direction for further work in attempting to understand these ESCGs as models of species coexistence and ecosystems biodiversity would be to generate results from similar RES experiments where the P -space path starts at a different initial point but still ends at (1,1,1), and also to run further RES-based simulations that end at (1,1,0) or (1,0,1), the other two P -space vertices where species extinctions can be expected to occur at the limit. The OES-based results can then be judged in the wider context of comparison to the array of results from the various RES-based experiments: the extent to which the sudden shift in $n_s(T)$ outcomes is also seen, if at all, in the RES experiments will then allow us to make a reasoned judgement on the generality of that shift occurring in OES-based outcomes. This is further work that I intend to report in one or more forthcoming papers, but at this stage it seems reasonable to hypothesize that the M -dependent shift in experiment outcomes seen in the OES systems as reported by (Zhong et al., 2022) and replicated here is in fact a mere artefact of the method, a quirk of the OES’s normalization and roulette-selection process.

6 CONCLUSION

The results presented here offer the first clear and comprehensive set of counterexamples to all the key results presented by (Zhong et al., 2022), because I have shown that in the RES-based ESCG experiments, which I contend are just as plausible models of real-world systems as the OES-based ones, no M -dependent shift in outcomes (i.e., final counts of coexisting species) occurs. Provision of these counterexamples, and my analysis and explanation in terms of P -space trajectories, are the contributions of this paper. I see no reason to expect that the problematic issues identified here with the use of OES in ESCGs is limited only to the one paper by Zhong et al. that I have focused on here. It seems plausible that con-

clusions drawn in many other papers reporting results from OES-based studies may need to be revised if, when the same M -sweep experiments are instead run with a RES-based ESCG, the results come out different. Only further work and time will tell, but right now, from what I've presented here, it does seem that the devil lies in the details for this class of evolutionary spatial ecosystem simulation model.

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