

LIC-R: Line Integral Convolution Revisited

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Abstract: We present a novel formulation of Line Integral Convolution (LIC), a fundamental method for visualizing vector fields in flow visualization. Our approach reinterprets the traditional LIC technique by leveraging a regularized, directional curvature flow along streamlines, utilizing material derivatives to achieve the desired convolution. By adopting an entirely Eulerian framework, our method eliminates the need for complex numerical integration and high-order interpolation schemes that are typically required in classical LIC algorithms. This shift not only simplifies the implementation of LIC, making it more accessible for both CPU and GPU architectures, but also significantly reduces the computational overhead. Despite these simplifications, our method maintains visual quality comparable to that of more traditional and computationally expensive approaches. Moreover, the discrete nature of our formulation makes it particularly well-suited for irregular grids and sparse data, broadening its applicability in practical settings. Through various experiments, we demonstrate that our algorithm delivers efficient and visually coherent results, offering an attractive alternative for dense flow visualization with reduced complexity.

1 INTRODUCTION

Flow visualization is a crucial area of scientific visualization focused on representing vector fields—such as fluid motion, air currents, or magnetic fields—in a way that makes their structure and dynamics visually understandable. By enabling the study of patterns like vortices, divergences, and flow separations, flow visualization is widely used in various fields, including fluid dynamics, meteorology, medical imaging, and aerodynamics (Delmarcelle and Hesslink, 2023). Accurate and efficient visualization of flow data is essential for understanding complex behaviors in natural and engineered systems, helping scientists and engineers in tasks ranging from analyzing weather patterns to optimizing the performance of mechanical components (Yau, 2024).

2023 marked the 30th anniversary of the Line Integral Convolution (LIC) method. Since its inception by (Cabral and Leedom, 1993), LIC has become one of the most popular and widely adopted techniques for dense and continuous visualization of vector fields. This method has been particularly valued

for its ability to generate high-quality visualizations of flow data. Over the past three decades, several improvements and variations of the LIC algorithm have been proposed to address its limitations and enhance its capabilities (Stalling and Hege, 1995; Laramée et al., 2008; Weinkauff and Theisel, 2010), solidifying its position as a go-to tool in flow visualization.

The classical LIC algorithm is conceptually simple, which is one of its key strengths. It begins with a random image $\mathcal{J} : \Omega \rightarrow [0, 1]$, where Ω is a d -dimensional domain, and a vector field $\mathcal{V} : \Omega \rightarrow \mathbb{R}^d$, defined over the same domain. The fundamental idea is to compute streamlines, which are curves that are everywhere tangential to the vector field \mathcal{V} , passing through each pixel of the image. Once streamlines are generated, the image \mathcal{J} is filtered along these streamlines using a low-pass convolution kernel. This process results in strong spatial coherence along the streamlines (i.e., along the flow direction) and reduced coherence across (i.e., perpendicular to the flow), producing the familiar streaked appearance characteristic of LIC visualizations. Traditionally, LIC is computed on a grayscale image, with color often reserved as an additional channel to encode other data properties. Figure 1 provides an illustrative example of such an LIC visualization.

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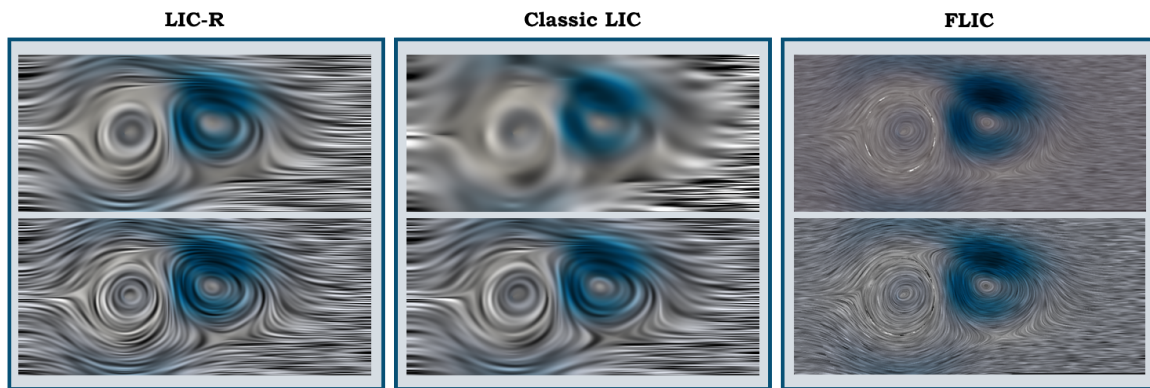


Figure 1: We present LIC-R, a novel efficient formulation for line integral convolution taking inspiration from the material derivative from continuum mechanics. Our method is able to capture the main features of flows with a quality comparable to classic LIC and fast LIC computed using high-order sampling and integration schemes but is significantly faster than classic LIC (up to one order of magnitude) and achieves better quality than Fast LIC with comparable processing times. In the example above, a frame from *jungelzbiemniak* simulation.

Despite its effectiveness, the classical LIC approach has several challenges and shortcomings. One significant issue is that the vector field is often available only at discrete sampling points, particularly when working with real-world or simulated data. To address this, LIC requires interpolation schemes to reconstruct the continuous vector field \mathfrak{V} from these discrete samples. The choice of interpolation can have a large impact on the quality of the resulting visualization. For high-quality LIC images, higher-order interpolation schemes, such as cubic Catmull-Rom splines (Catmull and Rom, 1974), are frequently employed. However, this introduces additional computational complexity.

Another challenge arises from the need to compute streamlines through numerical integration. The accuracy and smoothness of the streamlines are also highly dependent on the integrator used. High-order numerical integration schemes, such as the classical fourth-order Runge-Kutta method (Stalling, 1998), are often required for high-quality results. These higher-order methods, while effective, introduce significant computational overhead and complexity, both in terms of implementation and performance. Thus, while the original LIC algorithm is conceptually straightforward, its practical implementation can become quite intricate and computationally expensive due to the necessity of high-order interpolation and integration methods.

As a result, implementing classical LIC remains a complex and resource-intensive task rooted in the Lagrangian view of flow, which requires tracing individual particle paths through a continuous vector field. The reliance on this Lagrangian approach necessitates careful balancing of accuracy, com-

putational cost, and visual quality, often requiring trade-offs based on the specific application. Additionally, the need for high-order interpolation and integration makes the method less suited for real-time applications or situations where computational resources are limited.

Contributions. In this paper, we propose a novel approach to LIC that shifts away from the traditional Lagrangian view and adopts a fully Eulerian perspective. Rather than interpolating the vector field to reconstruct continuous flow paths, we work directly with the discrete vector samples. This shift allows us to reinterpret the LIC process as a discrete curvature flow along each streamline, using directional derivatives that naturally align with the parametric derivatives of traditional LIC. Our formulation not only simplifies the implementation but also significantly reduces computational costs by avoiding the need for high-order interpolation and numerical integration. Importantly, our method produces visual results that are comparable in quality to much more computationally expensive approaches.

Furthermore, since our technique does not rely on interpolation to a continuum, it is easily generalized to vector fields that are sampled on irregular grids, broadening the potential application of LIC to a wider range of data sources. This paper provides a detailed discussion of our novel formulation and demonstrates its effectiveness across several different test cases.

Organization. The remainder of this paper is structured as follows: In the next section, we review related work, focusing on major improvements and variations of the original LIC algorithm that have

been proposed over the past 31 years. Section 3 provides the mathematical foundations of our new Eulerian LIC approach, including the re-derivation of the LIC process using discrete curvature flow. Section 4 details the implementation of our algorithm and highlights its computational advantages. Finally, we present results and discuss limitations in Section 5, before concluding the paper with an outlook on future work.

2 RELATED WORK

Our method deals with image-based visualization of vector fields. Such fields have attracted the attention of the visualization community for a long time and we do not aim to provide a comprehensive discussion of the impressive body of relevant literature. Instead, we refer the reader to (Laramee et al., 2008) excellent survey on texture-based flow visualization as well as more recent surveys covering flow visualization for environmental science (Bujack and Middel, 2020), state-of-the-art methods for vortex extraction (Günther and Theisel, 2018), and time-dependent flow topology (Bujack et al., 2020).

Recent trends in flow visualization (Günther, 2020) are targeting the challenges related to automatic extraction and visual analysis of vortices and flow features by considering observer-relative frame references (Hadwiger et al., 2019; Rautek et al., 2021; Günther et al., 2017; Zhang et al., 2022; Günther and Theisel, 2024) or topology-based methods (Rojo and Günther, 2020; Hofmann and Sadlo, 2021). To this end, the community has started to explore modern machine learning, for example for performing neural flow map interpolation (Jakob et al., 2021), or for automatic extraction of reference frames through convolutional neural networks (Kim and Günther, 2019). Very recently, (Daßler and Günther, 2024) defined a theoretical framework for extracting flow features through variational methods.

In the following, we discuss the classical image-based methods most closely related to our formulation. Since the publication of (Cabral and Leedom, 1993) seminal paper, the visualization community directed massive efforts to improve the original formulation and to apply it to a variety of scientific domains. The original method basically involves the convolution of a noise texture with a kernel function along a characteristic path defined by the flow field in a way to produce a streaky visualization that highlights the structure of the flow and allows for immediate understanding of its features.

(Rezk-Salama et al., 1999) extended the formu-

lation to 2D surfaces through the usage 3D textures, (Stalling and Hege, 1995) derived a resolution-independent formulation that also addresses the original LIC’s high computational demands, (Wegenkittl et al., 1997) extended the original formulation to asymmetric kernels to generate oriented patterns that are able to highlight the flow orientation, (Forsell and Cohen, 1995) extended LIC to curvilinear grids, while (Sundquist, 2003) proposed a dynamic version for representing streamline evolution. LIC on surfaces is also discussed thoroughly in (Stalling, 1998) PhD thesis, as is the numerical machinery for that purpose. To address the computational cost of 3D LIC and surface LIC, image-based computation has been proposed (van Wijk, 2003; Telea and van Wijk, 2003) but the resulting visualization is not frame-coherent in animations.

Concerning the various applications, apart of direct visualization of flows for environmental sciences and engineering (Bujack and Middel, 2020), the method has been successfully applied in geometry processing for effective illustration of surface shapes (Interrante, 1997), for highlighting the salient regions in molecular surfaces (Lawonn et al., 2014), for terrain and cartography applications (Jenny, 2021), and for highlighting fiber tracts maps from diffusion tensor imaging (McGraw et al., 2002). Very recently, (Rautek et al., 2023) integrated line integral convolution in a framework for interactive analysis and identification of physically observable vortex structures. Inspired by advection and the concept of the material derivative in continuum mechanics, which refers to the rate of change over time of a physical quantity in a material element exposed to a velocity field (Sun, 2020), we propose a novel Eulerian formulation for line integral convolution. Our method is based on the material transport of underlying noise images. Eulerian formulations are popular for computing and rendering Finite-Time Lyapunov Exponents to characterize coherent Lagrangian structures (Garth et al., 2007; Sadlo and Peikert, 2007), or for performing post-hoc flow analysis (Agronovsky et al., 2014). (Hanser et al., 2019) recently also considered the material derivative concept for creating an energy-based visual analysis system for 2D flow fields.

3 METHODOLOGY

In this section, we detail our novel approach for computing Line Integral Convolution (LIC) using a regularized, directional curvature flow along streamlines in the vector field. Our method builds upon the clas-

sical LIC framework but simplifies the computation by eliminating the need for interpolation and numerical integration. Below is a brief outline of the key components:

- **Continuum Formulation.** We begin with the classical definition of LIC using streamlines and arc-length parameterization. Our method adapts this to directional derivatives along the vector field, allowing for efficient computation.
- **Material Derivative and Curvature Flow.** Instead of advection, we introduce a curvature-based flow derived from the second order directional derivative, analogous to the material derivative in fluid dynamics (Sun, 2020).
- **Discrete Setting.** The vector field and image are discretized on a grid, and we derive second and fourth order difference stencils for directional derivatives to efficiently compute the LIC result on both uniform and irregular grids.
- **Numerical Solution.** We present an iterative numerical solution using Gauss-Seidel steps for solving the resulting system matrix. The full solution is obtained by combining equilibrium and perturbed step solutions.

3.1 Continuum Formulation

Given a d -dimensional vector field $\mathfrak{V} : \Omega \rightarrow \mathbb{R}^d$ over a domain $\Omega \subseteq \mathbb{R}^d$, streamlines are defined as lines tangential to \mathfrak{V} . Given a point $\mathbf{x}_0 \in \Omega$, individual streamlines passing through \mathbf{x}_0 can be computed explicitly by forward and backward integration, resulting in positions $\mathbf{p}(\mathbf{x}_0, t)$ along a curve parameterized by t . From the definition it is clear that the resulting curve is Euler-invariant, that is, $\mathbf{p}(\mathbf{x}_0, t)$ does not change under different parameterizations. LIC uses this fact to extract curves under an arc-length parameterization s , $\mathbf{p}(\mathbf{x}_0, s)$, for which

$$\frac{d\mathbf{p}}{ds} = \frac{\mathfrak{V}(\mathbf{p}(s))}{\|\mathfrak{V}(\mathbf{p}(s))\|_2}. \quad (1)$$

LIC uses these streamlines to convolve a low-pass filter κ of length L with an image \mathcal{I} containing uncorrelated noise, to compute the final image \mathfrak{L} ,

$$\mathfrak{L}(\mathbf{x}_0) = \int_{-L/2}^{L/2} \kappa(s) \mathcal{I}(\mathbf{p}(\mathbf{x}_0, s)) ds. \quad (2)$$

Our approach, in contrast, uses Eqn. (1) to define *directional* derivatives:

$$\frac{d\mathbf{p}}{ds} = \frac{\partial \mathbf{p}}{\partial \mathbf{v}}, \quad (3)$$

where \mathbf{v} is a *normalized* (unit) vector of the underlying vector field. Using the chain rule, we have

$$\frac{\partial \mathbf{p}}{\partial \mathbf{v}} = \frac{\partial \mathbf{p}}{\partial \mathbf{x}} \frac{\partial \mathbf{x}}{\partial \mathbf{v}}. \quad (4)$$

The term $\partial \mathbf{x} / \partial \mathbf{v}$ in Eqn. (4) measures the change of \mathbf{x} in direction of \mathbf{v} , which is equivalent to the components of \mathbf{v} . We thus obtain

$$\frac{\partial \mathbf{p}}{\partial \mathbf{v}} = \langle \mathbf{v}, \nabla \mathbf{p} \rangle. \quad (5)$$

For the secondorder derivative, we analogously use

$$\frac{\partial^2 \mathbf{p}}{\partial \mathbf{v}^2} = \mathbf{v}^T H_{\mathbf{p}} \mathbf{v}, \quad (6)$$

where $H_{\mathbf{p}}$ is the Hessian of \mathbf{p} .

3.2 Material Derivative and Curvature Flow

Computing $\nabla \mathbf{p}$ and $H_{\mathbf{p}}$ seems problematic at first glance, since each individual path should be parametrized and the notion of a canonical coordinate frame \mathbf{x} changes along the path \mathbf{p} . It is at this point where we take inspiration from the material derivative. Given a material field $\mathcal{I}(\mathbf{x}, t)$ and a macroscopic velocity field (one that depends only on space-time) $\mathbf{v}(\mathbf{x}, t)$, the material derivative measures total change over time. It is defined as

$$\frac{D\mathcal{I}}{Dt} := \frac{\partial \mathcal{I}}{\partial t} + \langle \mathbf{v}, \nabla \mathcal{I} \rangle. \quad (7)$$

The two terms in Eqn. 7 model the change in material due to external ($\partial \mathcal{I} / \partial t$) and advection effects ($\langle \mathbf{v}, \nabla \mathcal{I} \rangle$), introducing an artificial integration time parameter t even for stationary fields. In our case, the first term vanishes. Since we are not interested in advecting the material but rather smoothing along each streamline, we define a *material curvature*, using the second order directional derivative instead, resulting in the following curvature flow.

$$\frac{D\mathcal{I}}{Dt} := \left| \frac{\partial^2 \mathcal{I}}{\partial \mathbf{v}^2} \right|. \quad (8)$$

This gives rise to a sequence of images $\mathcal{I}(\mathbf{x}, t)$ that is now decoupled from the underlying streamlines. Note that the 2nd order directional derivative is equivalent to a derivative with respect to arc-length and, thus, its modulus measures the curvature (Farin, 2001) of the image in direction \mathbf{v} . Note that in our use case, an arclength parameterization of characteristic curves can be obtained by pre-normalizing the vector field, except for stationary points which should remain at 0-magnitude (see also Sec. 4 on stationary points).

Table 1: Components of the central difference stencils (cf. (12)) used in this work. Adding the stencils in either the $O(\Delta\mathbf{x}^2)$ or $O(\Delta\mathbf{x}^4)$ columns yields 2nd or 4th-order discretisations of the second directional derivative with respect to $\mathbf{v} = (v_1, v_2)$. The velocity \mathbf{v} is sampled at the center of the stencil.

$\frac{\partial^2}{\partial \mathbf{v}^2}$	$O(\Delta\mathbf{x}^2)$	$O(\Delta\mathbf{x}^4)$
$\frac{\partial^2}{\partial x_1^2}$	$\frac{1}{\Delta x_1} \begin{bmatrix} 0 & 0 & 0 \\ 1 & -2 & 1 \\ 0 & 0 & 0 \end{bmatrix}$	$\frac{v_1}{12\Delta x_1^2} \begin{bmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ -1 & 16 & -30 & 16 & -1 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{bmatrix}$
$\frac{\partial^2}{\partial x_2^2}$	$\frac{1}{\Delta x_2} \begin{bmatrix} 0 & 1 & 0 \\ 0 & -2 & 0 \\ 0 & 1 & 0 \end{bmatrix}$	$\frac{v_1}{12\Delta x_2^2} \begin{bmatrix} 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 16 & 0 & 0 \\ 0 & 0 & -30 & 0 & 0 \\ 0 & 0 & 16 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 \end{bmatrix}$
$\frac{\partial^2}{\partial x_1 \partial x_2}$	$\frac{1}{4\Delta x_1 \Delta x_2} \begin{bmatrix} 1 & 0 & -1 \\ 0 & 0 & 0 \\ -1 & 0 & 1 \end{bmatrix}$	$\frac{v_1 v_2}{72\Delta x_1 \Delta x_2} \begin{bmatrix} 1 & -8 & 0 & 8 & -1 \\ -8 & 64 & 0 & -64 & 8 \\ 0 & 0 & 0 & 0 & 0 \\ 8 & -64 & 0 & 64 & -8 \\ -1 & 8 & 0 & -8 & 1 \end{bmatrix}$

3.3 Discrete Setting

For the sake of clarity, we assume for now that both the vector field \mathfrak{V} and the image \mathfrak{I} are sampled on a uniform 2D grid. However, neither dimension nor uniformity of the grid are a requirement of our method. To discretize Eqn. (8) we use discrete differences. We consider both second order and fourth order stencils, which can be derived in the classical fashion (i.e., using a Taylor expansion around any grid point). Expanding the second order derivative in Eqn. (8) using the formulation in Eqn. (6), we obtain

$$\frac{\partial^2 \mathfrak{I}}{\partial \mathbf{v}} = v_1^2 \frac{\partial^2 \mathfrak{I}}{\partial x_1^2} + v_2^2 \frac{\partial^2 \mathfrak{I}}{\partial x_2^2} + 2v_1 v_2 \frac{\partial^2 \mathfrak{I}}{\partial x_1 \partial x_2}, \quad (9)$$

where subscripts denote vector components. The second order discrete difference stencils are

$$\begin{aligned} \frac{\partial^2}{\partial x^2} &= \frac{1}{\Delta x^2} [1 \quad -2 \quad 1] + O(\Delta x^2) \text{ and} \\ \frac{\partial}{\partial x} &= \frac{1}{2\Delta x} [-1 \quad 0 \quad 1] + O(\Delta x^2). \end{aligned} \quad (10)$$

The fourth order counterparts are

$$\begin{aligned} \frac{\partial^2}{\partial x^2} &= \frac{1}{12\Delta x^2} [-1 \quad 16 \quad -30 \quad 16 \quad -1] + O(\Delta x^4) \\ \text{and} \end{aligned} \quad (11)$$

$$\frac{\partial}{\partial x_1} = \frac{1}{12\Delta x} [1 \quad -8 \quad 0 \quad 8 \quad -1] + O(\Delta x^4).$$

Combining these stencils along directions x_1, x_2 and assuming unit Δx , we assemble 2D stencils (3×3

and 5×5 , respectively) for the second order directional derivative (\otimes denotes an exterior vector product, interpreting vectors as matrices and giving rise to a matrix),

$$\frac{\partial^2}{\partial \mathbf{v}^2} = v_1^2 \frac{\partial^2}{\partial x_1^2} + v_2^2 \left(\frac{\partial^2}{\partial x_2^2} \right)^T + 2v_1 v_2 \frac{\partial}{\partial x_1} \otimes \left(\frac{\partial}{\partial x_2} \right)^T, \quad (12)$$

that takes a form of a matrix with circulant-like structure, that in the following we will refer to as

$$\mathcal{D} := \frac{\partial^2}{\partial \mathbf{v}^2}. \quad (13)$$

In this way, the discrete problem takes the following form.

$$\frac{D\mathfrak{I}}{Dt} = \mathcal{D} \star \mathfrak{I}. \quad (14)$$

Table 1 summarizes the components of the second directional derivative stencils used in this work. Summing up the stencil contributions of either the 2nd or 4th order stencil yields the final discrete difference operator, \mathcal{D} . The directional derivative estimate of a noise image \mathfrak{I} with respect to the vector field \mathbf{v} is then obtained by computing the non-homogeneous convolution, $\mathcal{D} \star \mathfrak{I} \approx \frac{\partial^2}{\partial \mathbf{v}^2} \mathfrak{I}$.

3.4 Numerical Solution

To solve the problem in Eqn. 14, we consider the superposition principle, and we separate it in two different contributions:

- a solution at equilibrium, obtained when the time derivative $D\mathcal{J}/Dt$ vanishes;
- a perturbed step solution, obtained through a regularized Euler integration step.

Our approach assembles the system matrix of the above expression on-the-fly and solves iteratively using Gauss-Seidel steps.

Equilibrium Solution. The equilibrium contribution \mathcal{J}_{eq} is the solution of the equation

$$\mathcal{D} \star \mathcal{J}_{\text{eq}} = 0. \quad (15)$$

Equation 15 essentially boils down to replacing the center pixel under the stencil \mathcal{D} with a weighted average of its neighboring pixels. The above condition is met if the weighted center pixel intensity, $c\mathcal{J}_{\text{eq}}(i, j)$ equals the negative sum, $-S_{i,j}$, of its neighbors (cf. Eqn. 16, below). In what follows, we are going to break up the above convolution to derive our iterative solver.

Since the matrix \mathcal{D} possesses a circulant-like structure that contains derivative operators, it is by definition singular. To solve it, we use an iterative Gauss-Seidel scheme together with Tikhonov regularization (Calvetti et al., 2000) and Jacobi preconditioning (Chow et al., 2018). For applying Jacobi method, we split the convolution by separating the diagonal contribution, which allows us to rewrite Eqn. 15 in the following form.

$$c(v_1^2 + v_2^2 + \mu)\mathcal{J}_{\text{eq}}(i, j) + S_{ij} = 0, \quad (16)$$

where c is the central co-factor of the stencil (either -2 or -30/12), μ is a small Tikhonov regularizer threshold for avoiding numerical instabilities, and S_{ij} is the rest of the stencil, *excluding* its center, convolved with the neighborhood of $\mathcal{J}_{\text{eq}}(i, j)$. Eqn. 16 can be used for iterative computation of image values $\mathcal{J}_{\text{eq}}(i, j)$ in the following way.

$$\mathcal{J}_{\text{eq}}^{(n)}(i, j) = \frac{-S_{ij}^{(n)}}{c(v_1^2 + v_2^2 + \mu)}, \quad (17)$$

where n is an iteration counter. In all results generated for this paper, we used $\mu = 10^{-6}$.

Perturbed Step Solution. For obtaining the full perturbed solution \mathcal{J} , we plug the equilibrium solution on a regularized forward Euler integrator, considering the following update equation:

$$\mathcal{J}^{(n+1)}(i, j) = \beta\mathcal{J}^{(n)}(i, j) + (1 - \beta)\mathcal{J}_{\text{eq}}^{(n)}(i, j). \quad (18)$$

For both second and fourth order (Eqn. (10), Eqn. (11)) derivative filters, we use $\beta = 0.2$, respectively.

4 IMPLEMENTATION DETAILS

In this section, we share some of our implementation details. We compare our proposed formulation (marked LIC-R in the rest of the manuscript), with the classical LIC formulation (marked LIC), and the fast LIC formulation (marked FLIC). We also tested our formulation against noise images generated with different distributions.

Numerical details. For numerical stability reasons, in the current implementation we consider a normalized vector field scenario. To this end, we first normalize every vector in the vector field except for $\mathbf{0}$. We then use a ping-pong buffer scheme (one read-only buffer for t and one write-only buffer for $t + 1$) to mimic Gauss-Seidel steps to iteratively solve Eqn. (18). After each iteration, we rescale the image to the original $[0, 1]$ range to mitigate amplification effects. The parameter β acts as an under-relaxation parameter in the iterative solver. In our experiment, directional diffusion did not produce meaningful results, underscoring the importance of the “cross-derivative” term $\partial^2\mathcal{J}/(\partial x_1\partial x_2)$. We attribute this to the geometric interpretation of $|v^T H_{\mathcal{J}} v|$ as directional curvature.

Stationary Points. In our current implementation, we implicitly assume that all vectors in the vector field have unit length. In real-world scenarios, this is not the case. In particular, Eqn. (15) is invalid at locations \mathbf{x} with $\mathfrak{V}(\mathbf{x}) = 0$. However, since our approach only considers samples at grid locations, such stationary points can be treated by skipping the iterative solver for these grid locations. Alternatively, one could argue that stationary points, especially those inside grid cells, are not well captured by LIC anyway and choose to ignore the issue. It is worth noting that either decision using our method leads to a more graceful behavior than the original LIC algorithm: If, in the original LIC algorithm, the streamline is not terminated at stationary points, the noise value in the scalar image may receive a relatively high weight due to the accumulation of multiple kernel coefficients, leading to artifacts.

Classic LIC. Our implementation of the classic LIC algorithm is based on regular (non-embedded) Runge-Kutta solvers of orders one through four (Stalling, 1998) paired with a choice of linear interpolation, C^1 -continuous Catmull-Rom interpolation (Catmull and Rom, 1974), and a Catmull-Rom style, C^2 -continuous quintic interpolation scheme. Since we observed diminishing returns

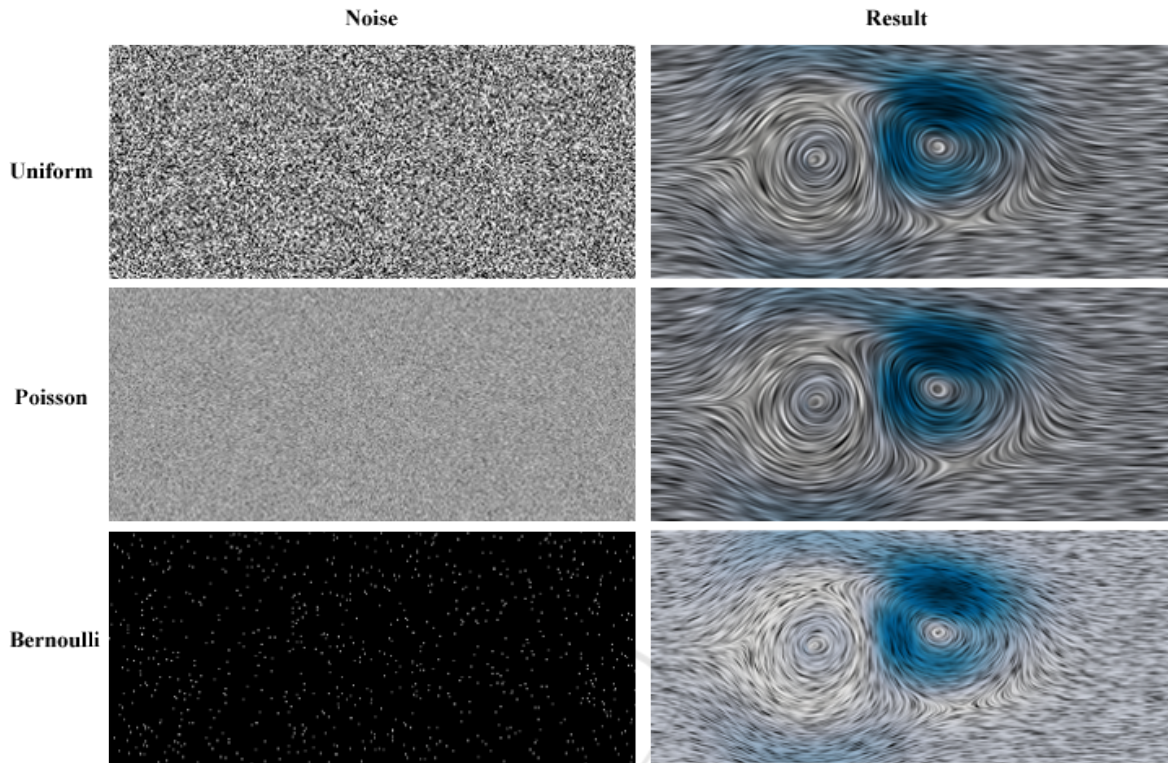


Figure 2: **Noise distribution comparison.** Our formulation is able to provide consistent results independently from the random noise image considered. In the figure a comparison between uniform, Poisson and Bernoulli noise.

Table 2: Comparison between classical LIC with various interpolators/integrators, FLIC, and LIC-R (average processing time in seconds).

Dataset	LIC			FLIC			LIC-R
	Lin./Euler	Lin./RK4	Cub./Euler	Lin./Euler	Lin./RK4	Cub./Euler	Order 4
jungtelziemniak2d	0.657	1.441	1.677	0.124	0.269	0.268	0.129
beads2d	0.105	0.239	0.284	0.012	0.017	0.017	0.052
boussinesq2d	0.606	1.405	1.319	0.138	0.154	0.173	0.150
cylinder2d	0.484	1.196	1.428	0.070	0.253	0.251	0.115
doublegyre2d	0.311	0.723	0.854	0.042	0.088	0.088	0.089
forcedampedduffing2d	0.105	0.241	0.294	0.011	0.016	0.016	0.025
fourcenters2d	0.116	0.369	0.284	0.012	0.020	0.020	0.043
pipedcylinder2d	0.504	1.029	1.222	0.168	0.202	0.200	0.104

between the C^1 and C^2 interpolators, we did not derive a “septic” C^3 -continuous version that would be the ideal pairing for the classical 4th order Runge-Kutta scheme. For both classic LIC and Fast LIC, we re-normalize vectors after interpolation to approximate spherical interpolation and conserve arc-length parameterization.

Fast LIC. Our implementation of the Fast LIC algorithm leverages the same integrators and samplers used for classic LIC, while exploiting streamline redundancy that occurs in LIC as a new streamline is always advected for each pixel despite the fact that many pixels are hit by the same streamline (Stalling and Hege, 1995). We consider equidistant streamline

sampling in combination with a box kernel to convert texture convolution to a texture-averaging operation. Thus, a streamline obtained to determine a pixel value can be reused with only minor corrections (due to the bi-directional kernel offset along the streamline) since two immediate neighboring pixels have nearly the same correlated pixels as their texture contributors. We maintain a pixel hit buffer to evaluate streamline coverage in the flow field, and therefore, only a small number of streamlines are actually advected for all pixels to collect texture contributions.

Noise Distributions. We tested our method on a variety of sparse noise patterns, generated through the Mersenne Twister random number generator (Mat-

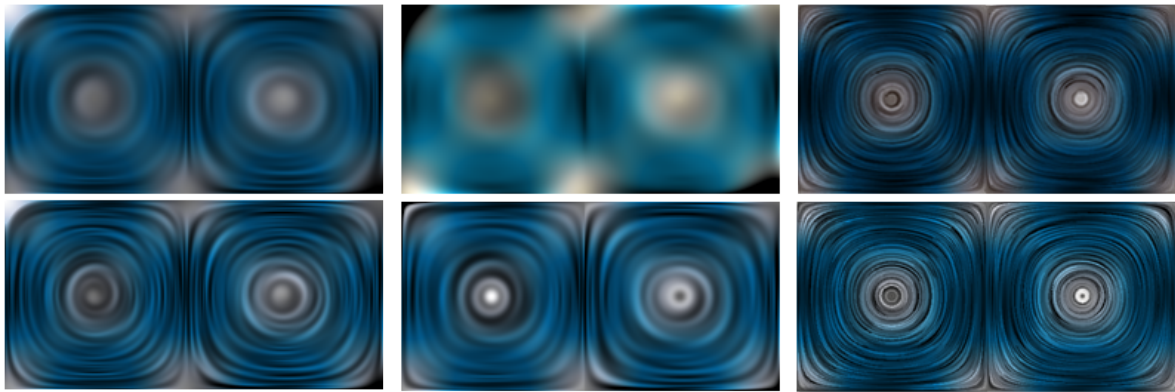


Figure 3: **Double gyre comparison.** From left to right, our LIC-R formulation, classic LIC, and Fast LIC. From top to bottom, for LIC-R: 2nd and 4th order derivative stencil; for Classic LIC and Fast LIC: linear and cubic with Runge-Kutta 4th order integration. Note that Fast LIC has a crisper appearance due to the limited number of integration steps, while LIC-R provides a smoother appearance.

sumoto and Nishimura, 1998) according to the following distributions:

- **Uniform.** We use the standard uniform generator that produces random floating-point values x , uniformly distributed on the interval $[a, b)$, that is, distributed according to the probability density function (Goulard, 2022):

$$P(x|a, b) = \frac{1}{b - a}.$$

- **Poisson.** We use the standard Poisson generator that produces random non-negative integer values i , distributed according to discrete probability function (Inouye et al., 2017):

$$P(i|\mu) = \frac{e^{-\mu} \mu^i}{i!}$$

The value obtained is the probability of exactly i occurrences of a random event if the expected, *mean* number of its occurrence under the same conditions (on the same time/space interval) is μ .

- **Bernoulli.** We use the standard Bernoulli generator that produces random boolean values, according to the discrete probability function (Dai et al., 2013). The probability of *true* is

$$P(b|p) = \begin{cases} p, & \text{if } b \text{ is true} \\ 1 - p, & \text{if } b \text{ is false} \end{cases}$$

In Fig. 2 we show examples of our method applied to a vector field with different noise distributions. The results showcased in this manuscript are generated with uniform noised distribution.

5 RESULTS

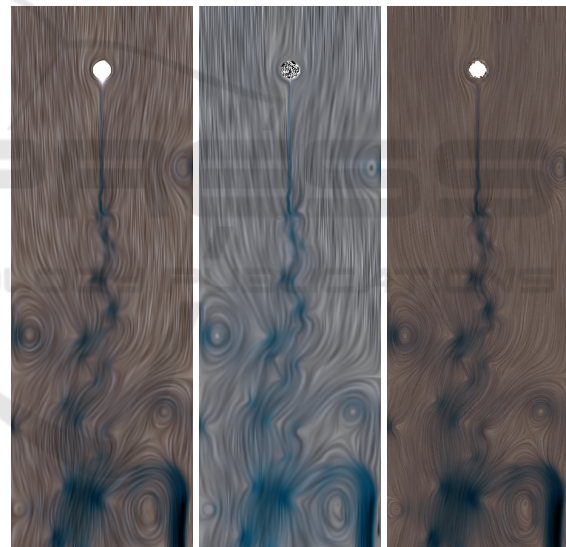


Figure 4: **Boussinesq comparison:** Left: LIC-R 4th order. Center: Classic LIC with cubic interpolator and Runge-Kutta 4th integrator. Right: Fast LIC with cubic interpolator and Runge-Kutta 4th integrator.

We benchmarked our method against the classical LIC (Cabral and Leedom, 1993) and the fast LIC (Stalling and Hege, 1995) on a variety of commonly-used 2D datasets: the simulation of a flow around a cylinder (von-Karman Vortex Street, “cylinder”), a simulation of a viscous 2D flow around two cylinders (“piped”), the simulation of a 2D flow generated by a heated cylinder in which the Boussinesq approximation is used (“boussinesq2d”), an analytical periodic time-dependent vector field, in which a

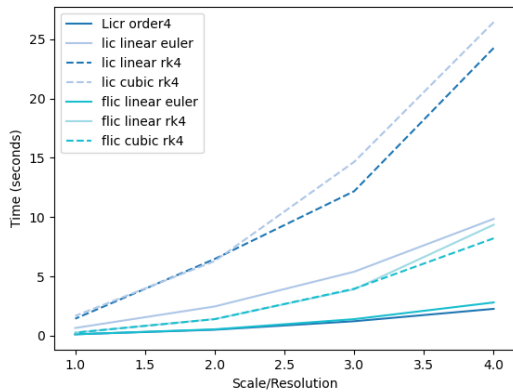


Figure 5: **Execution time comparison:** The plot shows the performance of several numerical techniques. The x-axis represents the resolution scale, while the y-axis reflects the computation time in seconds. The result highlights LIC-R 4th order maintains execution times comparable to Fast LIC across the scales.

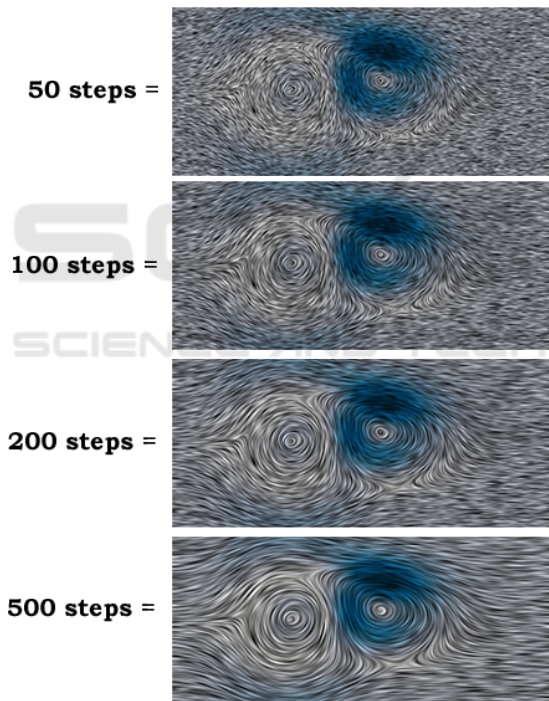


Figure 6: **LIC-R convergence:** From top to bottom, the LIC-R 4th order for a varying number of Gauss-Seidel steps.

separating boundary oscillates horizontally between two oppositely rotating vortices (“double-gyre”), an analytic data set of a rotating Petri dish, in which the vortex center moves on a circular path (Weinkauff and Theisel, 2010) (“petri”), an analytic data set containing four vortices (“four”), an analytic vector field describing the phase space of a duffing oscillator with forcing and damping (“duffing”), a synthetic vector field representing a von-Karman vortex street genera-

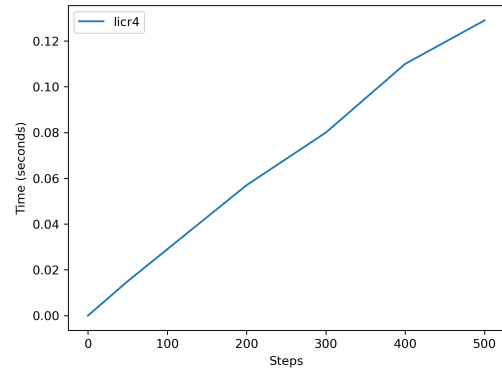


Figure 7: **LIC-R complexity:** The graph shows the computation time for LIC-R 4th order at the varying of integration steps, showcasing the linear complexity.

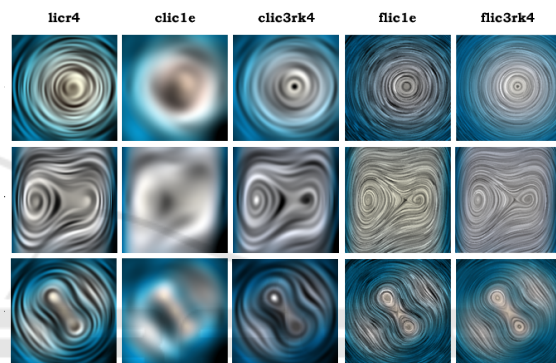


Figure 8: **Qualitative results:** From left to right, comparison of our method (LIC-R order 4) with Classic LIC (linear plus Euler, and cubic plus Runge-Kutta 4th order), and Fast LIC (linear plus Euler, and cubic plus Runge-Kutta 4th order). From top to bottom, *beads2d*, *forceddampedduffing*, and *fourcenters*. Note that Fast LIC has a crisper appearance due to the limited number of integration steps, while LIC-R provides a smoother appearance.

tion and constructed as co-gradient to a stream function (“jungtel”). The datasets can be downloaded from ETH Zurich’s repository¹. The tests were performed on a Razer Blade Stealth laptop equipped with a CPU Intel i7-8565U, 4 cores at 1.8 GHz. For visualization purposes, we show the vector field modulus/magnitude color-mapped through ColorBrewer’s colormap PuBu (Purple-Blue). For comparison, we considered the classic and the Fast LIC formulation with a variety of samplers and integrators: namely, the samplers range from linear up to cubic order, while the integrators range from the simple Euler up to the 4th order Runge-Kutta. Table 2 shows the average processing times in seconds for our method compared to various versions of LIC: our formulation is three times faster than the simplest LIC implementation (linear sampler with Euler integrator),

¹<https://cgl.ethz.ch/research/visualization/data.php>

and one order of magnitude faster than more complex LIC formulations (cubic sampling through Catmull-Rom splines and integration through 4th order Runge-Kutta), while at the same time is maintaining computation time comparable to Fast LIC implementation. Figure 5 compares LIC-R execution times for the *jungtelziemniak2d* dataset with classic LIC and fast LIC across different image scaling factors. It showcases that computation times are comparable to fast LIC, which involves a significant implementation overhead due to caching of partial characteristic lines. Figure 7 shows the computation times for the *jungtelziemniak2d* dataset obtained with different iteration steps, showcasing that our method has linear complexity with respect to the number of steps. Figure 6 shows the output of LIC-R 4th order for different integration steps. Our examples also corroborate that higher order integrators need to be paired with higher order interpolators, as shown in Fig. 3, Fig. 4 and Fig. 8, and as highlighted in (Stalling, 1998) thesis. In these figures, we provide qualitative comparisons between the classic LIC with different samplers and integrators and our method. We use a 3-tap tent filter and 500×1 pixel iterations for all classic LIC figures in this paper, and 500 iterations with $\beta = 0.2$ for LIC-R, while we use a 51-tap triangular kernel for Fast LIC and 5 iterations. We believe it is evident that LIC-R provides a quality comparable to the original method using Catmull-Rom interpolation and 4th order Runge-Kutta integration. In particular, we observe that classic LIC tends to blur the resulting image earlier, which we conjecture is a result of accumulating more numerical round-off errors, while fast LIC provides crisper images but the coherent trajectories seem to be shorter when compared to LIC-R. This can arguably be remedied by using bigger kernel sizes but this choice comes at a computational penalty. The accompanying video provides additional examples.

Limitations. We already highlighted the role of stationary points but would argue that such points are better added to the visualization using additional methodology that is orthogonal to LIC. Another limitation of our method is that, in the current implementation, we cannot control the convolution kernel κ , the kernel arises naturally as the solution to the curvature flow problem. A minor limitation is the non-isometric version of the derivative stencils used in this work: our stencils only have a 90° rotational invariance, yet previous work shows that using stencils with higher rotational invariances (or even isometric stencils) can result in substantially better results (Kamgar-Parsi et al., 1999). In our method, the construction of isometric differential stencils is hindered, but we

believe it is not made impossible by the use of directional derivatives.

6 CONCLUSIONS

We have presented a novel formulation of the original LIC algorithm that takes inspiration of material derivatives: material (the noise image) is smoothed out along streamlines using curvature flow. Unlike previous methods, we take a fully Eulerian view of the problem, resulting in significant improvements in terms of speed and ease-of-implementation over the classical formulation.

In the future, we plan to improve the isometry of our derivative stencils and investigate how custom smoothing kernels can be integrated into the framework. We also plan to apply the method to 2-manifolds embedded in 3D space by combining our method with approaches such as the Closest Point Method (MacDonald and Ruuth, 2008; Auer et al., 2012) for solving partial differential equations. A GPU-based implementation is likewise on our agenda, as is a closer look at the material derivative formulation for unsteady flows. Note that our method only has one pre-requisite, the formulation of a directional, second-order derivative, which is well-studied on 2-manifolds in the context of curvature. However, sampling the noise image \mathcal{J} in the vicinity of the 2-manifold is non-trivial, and more research with a special attention to computational demands is needed.

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